

EWK precision measurements

[present and future]

W. Hollik
14/4/99

Higgs: only unknown of SM

Higgs and **top** affect relations between high precision EW observables, through **quantum effects**

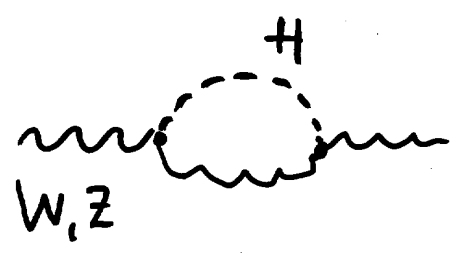
$$O_i = f_i(M_W, M_Z, M_H, m_t, \dots)$$

→ constraints on M_H

confronted with direct measurements

- tests of consistency
- signals for "new physics"

Higgs → precision observables



Higgs - gauge coupling

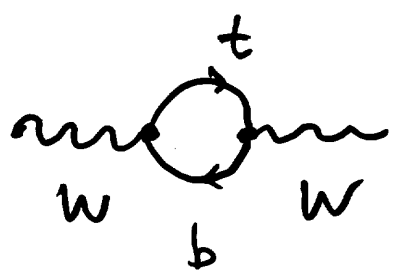
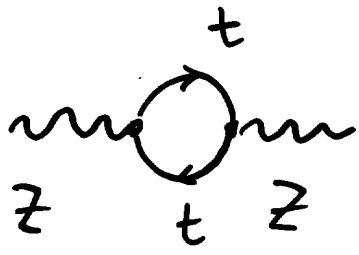
at 2-loop level:

Higgs - Yukawa interaction

Higgs - self interaction

top \rightarrow precision observables

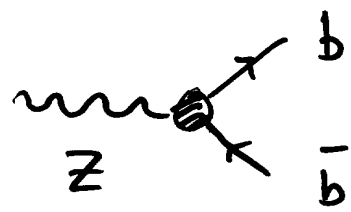
(1)



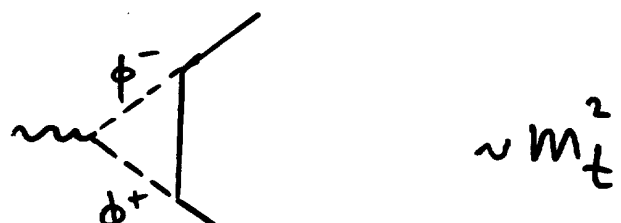
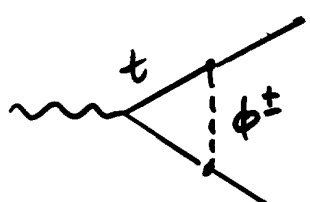
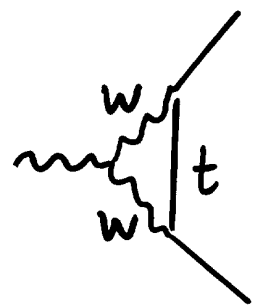
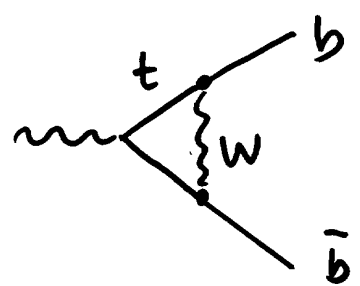
\Rightarrow ρ -parameter, Δr , g_V^f, g_A^f

$\left[\begin{array}{c} \text{Feynman diagram: top quark loop with photon external lines} \\ \gamma \quad t \quad \gamma \end{array} \right] \sim Q_t^2 \frac{M_Z^2}{m_t^2} \text{ (small)}$

(2)



$\Rightarrow R_b [A_b, A_{FB}^b]$



$\sim m_t^2$

Electroweak parameters⁴

$$\bullet \quad G_\mu = \frac{\pi}{\sqrt{2}} \cdot \frac{\alpha}{M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right)} \underbrace{\left[1 + \Delta\alpha - \frac{C_W^2}{S_W^2} \Delta\mathcal{P} + \dots \right]}_{= \Delta r}$$

$$S_W^2 = 1 - \frac{M_W^2}{M_Z^2}$$

$$C_W^2 = \frac{M_W^2}{M_Z^2}$$

$$\bullet \quad \sin^2 \theta_{\text{eff}}^l(M_Z) = S_W^2 + C_W^2 \Delta\mathcal{P} + \dots$$

• NC coupling (Zff), overall normalization

$$\left(\sqrt{2} M_Z^2 G_\mu \right)^{1/2} \left[1 + \frac{1}{2} \Delta\mathcal{P} + \dots \right]$$

Vector boson masses

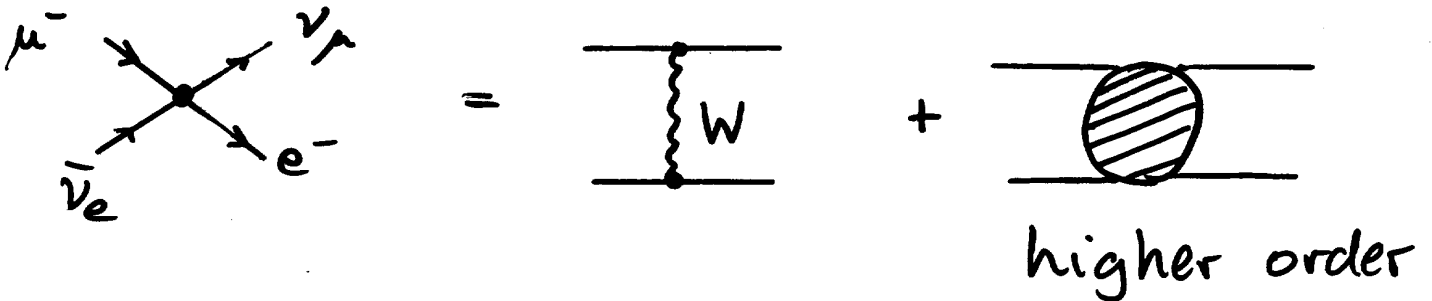
correlated by μ -lifetime $\leftrightarrow G_\mu$

$$\frac{1}{\tau_\mu} = \frac{G_\mu^2 m_\mu^5}{192 \pi^3} \left(1 - \frac{8m_e^2}{m_\mu^2}\right) \left\{ 1 - 1.81 \frac{\alpha}{\pi} + 6.70 \left(\frac{\alpha}{\pi}\right)^2 \right\}$$

↑
Sirlin
et al.

↑
van Ritbergen
Stuart

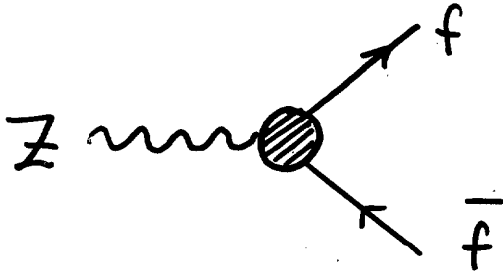
$$G_\mu = 1.16637(1) \cdot 10^{-5} \text{ GeV}^{-2}$$



$$G_\mu = \frac{\pi \alpha}{\sqrt{2}} \cdot \frac{1}{M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right)} \left(1 + \Delta r + \dots\right)$$

$$\Delta r = \underbrace{\Delta \alpha}_{\uparrow 6\%} - \underbrace{\frac{c_W^2}{s_W^2} \Delta \rho}_{\sim 4\%} + (\Delta r)_{\text{rem}} \left(m_t, M_H\right) \sim 1\%$$

resum: $\alpha (1 + \Delta \alpha + \dots) \rightarrow \frac{\alpha}{1 - \Delta \alpha} = \alpha(M_Z)$ LL-resu

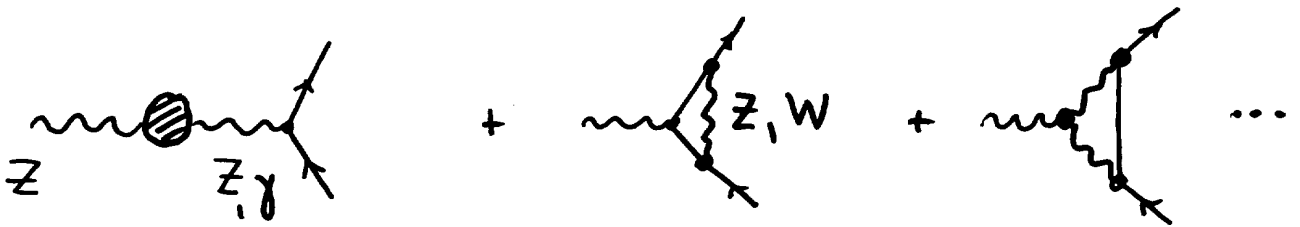


$$g_A^f = \sqrt{\rho_f} I_3^f$$

$$g_V^f = \sqrt{\rho_f} (I_3^f - 2 Q_f \sin^2 \theta_f)$$

$$\rho_f(m_t, M_H, \alpha_s) = \frac{1}{1 - \Delta\rho} + \dots$$

$$\sin^2 \theta_f(m_t, M_H, \alpha_s) = s_W^2 + \Delta\rho \cdot c_W^2 + \dots$$



non-universal

higher order EW (2-loop):

$$A G_\mu^2 m_t^4 + B G_\mu^2 M_Z^2 m_t^2 + (\text{unknown})$$

\uparrow
 $\Delta\rho \quad \Delta\rho^2$

\uparrow
Degrassi et al.

(i)

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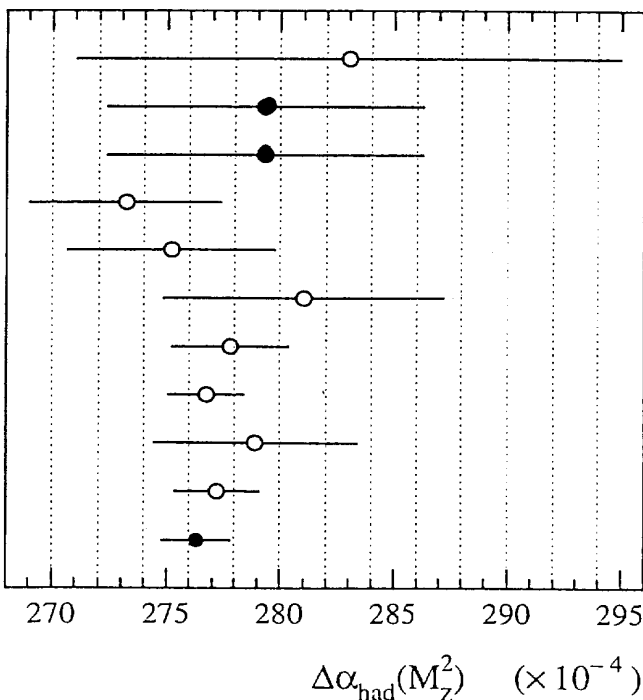
$$\Delta\alpha = (\Delta\alpha)_{lep} + (\Delta\alpha)_{had}^{(5)} + (\Delta\alpha)_{top}$$

$$(\Delta\alpha)_{lept} = \sum_{\gamma}^{e, \mu, \tau} m_{\gamma} \text{O} m_{\gamma} + m_{\gamma} \text{O} m_{\gamma} + m_{\gamma} \text{O} m_{\gamma}$$

Källén, Sabry (1955) Steinhauser (1998)

$$(\Delta\alpha)_{had}^{(5)} = -\frac{M_Z^2}{4\pi^2\alpha} \text{Re} \int_{4m_{\pi}^2}^{\infty} ds \frac{\sigma(e^+e^- \rightarrow had)}{s - M_Z^2 - i\epsilon}$$

$$\alpha(M_Z) = \frac{\alpha}{1 - \Delta\alpha} = \alpha [1 + \Delta\alpha + \Delta\alpha^2 + \dots]$$



- Lynn, Penso, Verzegnassi, '87
- Eidelman, Jegerlehner '95
- Burkhardt, Pietrzyk '95
- Martin, Zeppenfeld '95
- Swartz '96
- Aleman, Davier, Höcker '97
- Davier, Höcker '97
- Kühn, Steinhauser '98
- Groote et al. '98
- Erlar '98
- Davier, Höcker '98

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T: "theory driven"

$$(\Delta\alpha)_{\text{lept}} = 0.0315$$

$$(\Delta\alpha)_{\text{had}} = 0.0280 \pm 0.0007 \quad \begin{array}{l} \text{Eidelman,} \\ \text{Fegerlehner} \\ \text{('95)} \end{array}$$

$$0.027782 \pm 0.000254$$

Fegerlehner ('99)

$$0.02763 \pm 0.00016$$

Davier, Höcker ('9)

$$0.02777 \pm 0.00017$$

Kühn, Steinhauser
(198)

(ii) ρ -parameter (heavy top)

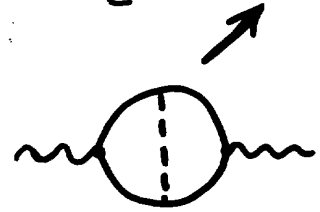
$$m_t \gg m_b$$

$$(\Delta\rho)_{1\text{-loop}}^{\text{top}} = 3 X_t$$

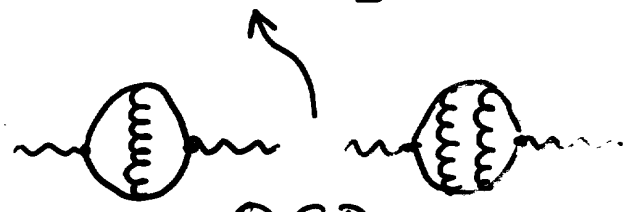
$$X_t = \frac{G_\mu m_t^2}{8\pi^2 \sqrt{2}} \quad \text{Veltman}$$

Higher order:

$$\Delta\rho = 3 X_t \left[1 + g^{(2)} X_t + \delta\rho_{\text{QCD}} \right]$$



2-loop EW



QCD

2- and 3-loop

van der Bij, Heugtenen

Bambieri et al.

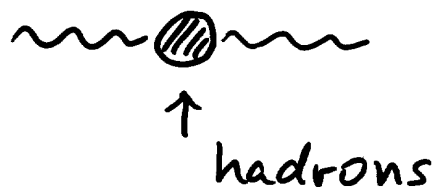
Fleischer, Jegerlehner, Tausov

{ Djouadi, Nguyen-Manh, Kniehl, Sirlin, Bardin et al., Djouadi, Gambino (2-loop)

{ Avdeev, Fleischer, Mikhailev, Tausov, Chetyrkin, Kühn, Steinhauser (3-loop)

Uncertainties of predictions

- $\delta(\Delta\alpha)$



$$\delta \sin^2 \theta = 0.00024$$



$$5 \cdot 10^{-5}$$

Eidelman,
Fegerlehner

Davies, Höcker
Kühn, Steinhauser

- QCD uncertainty

$\delta\alpha_s$ exp.
 scale uncertainty
 m_b uncertainty

} small

$$\delta \sin^2 \theta_e \approx 3 \cdot 10^{-5}$$

$$\delta(\Delta p) \approx 1 \cdot 10^{-4}$$

- uncertainties from unknown higher order EW contributions

- different schemes
- different options for

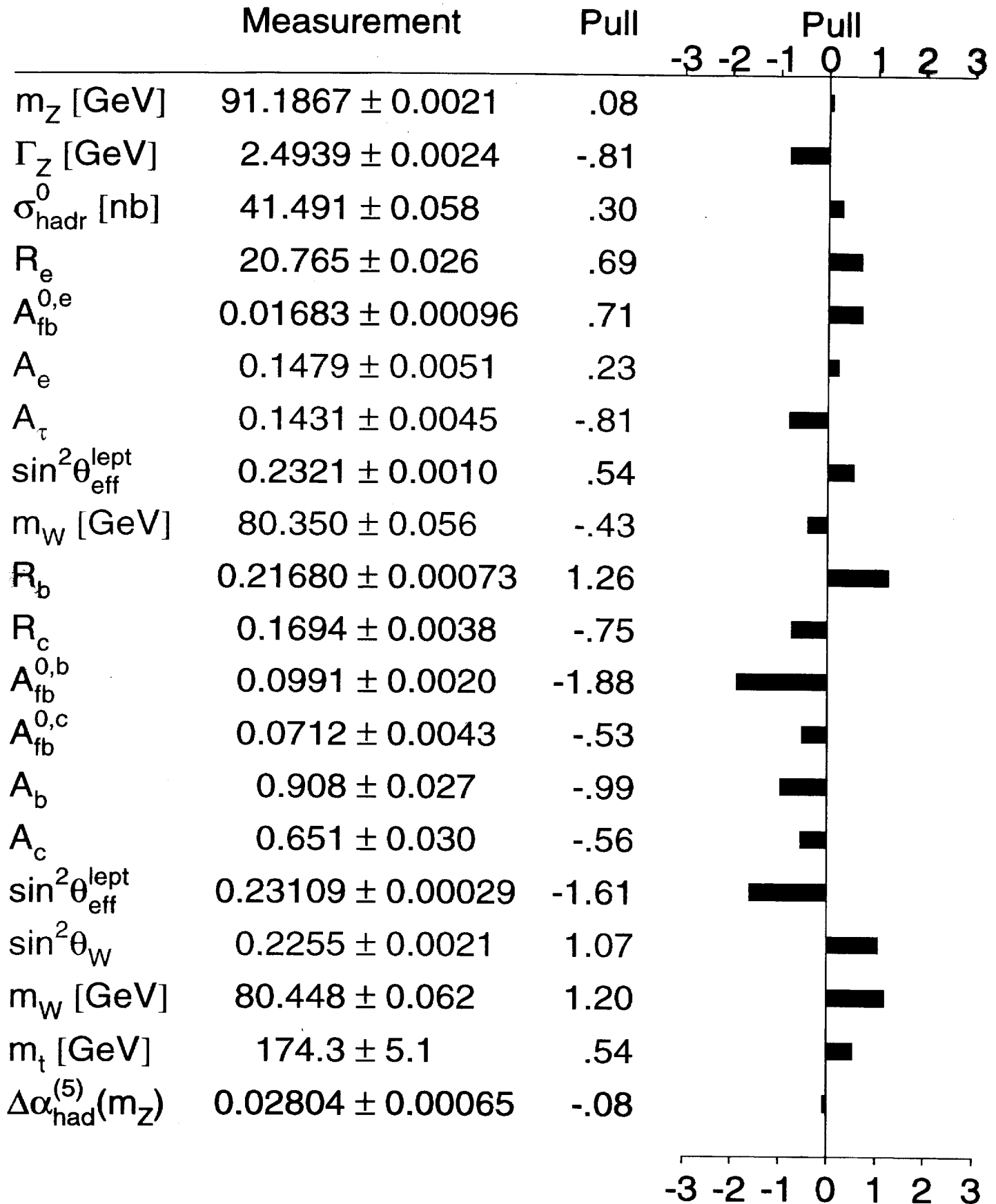
uncontrolled higher order terms

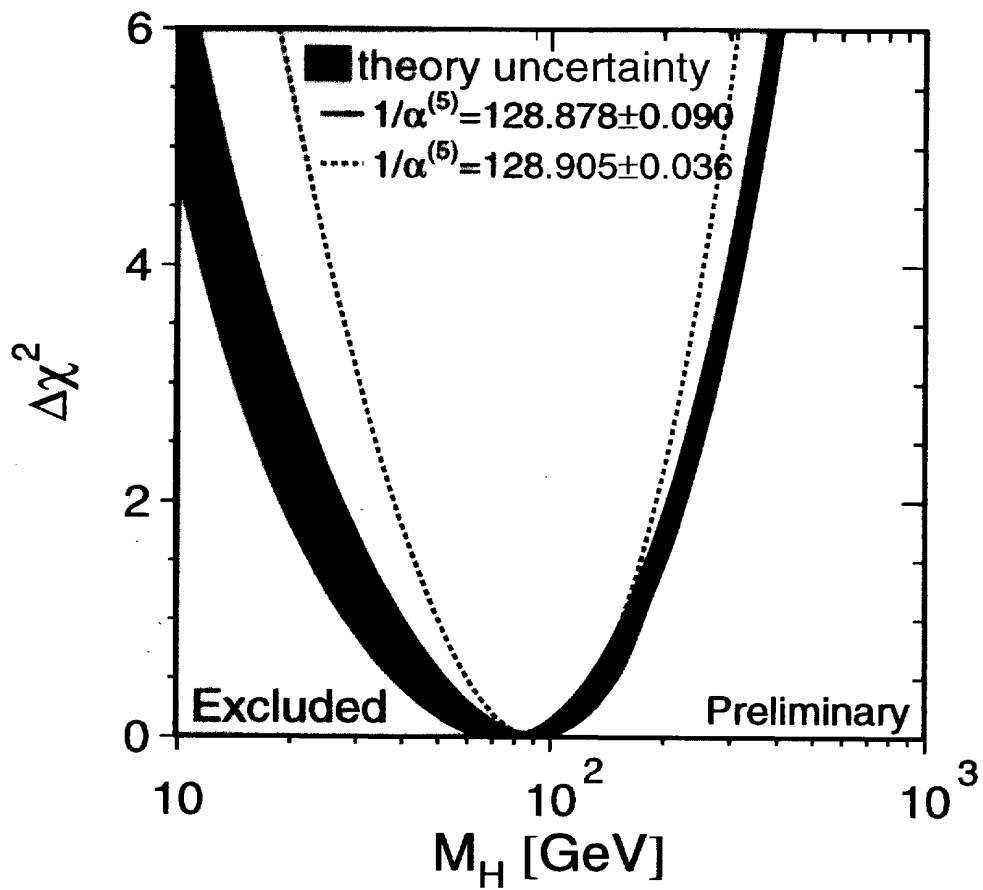
CERN 95-03 (eds. Bardin, Witt, Passarino)

Bardin, Grünewald, Passarino, hep-ph/9702452

MORIOND 1999

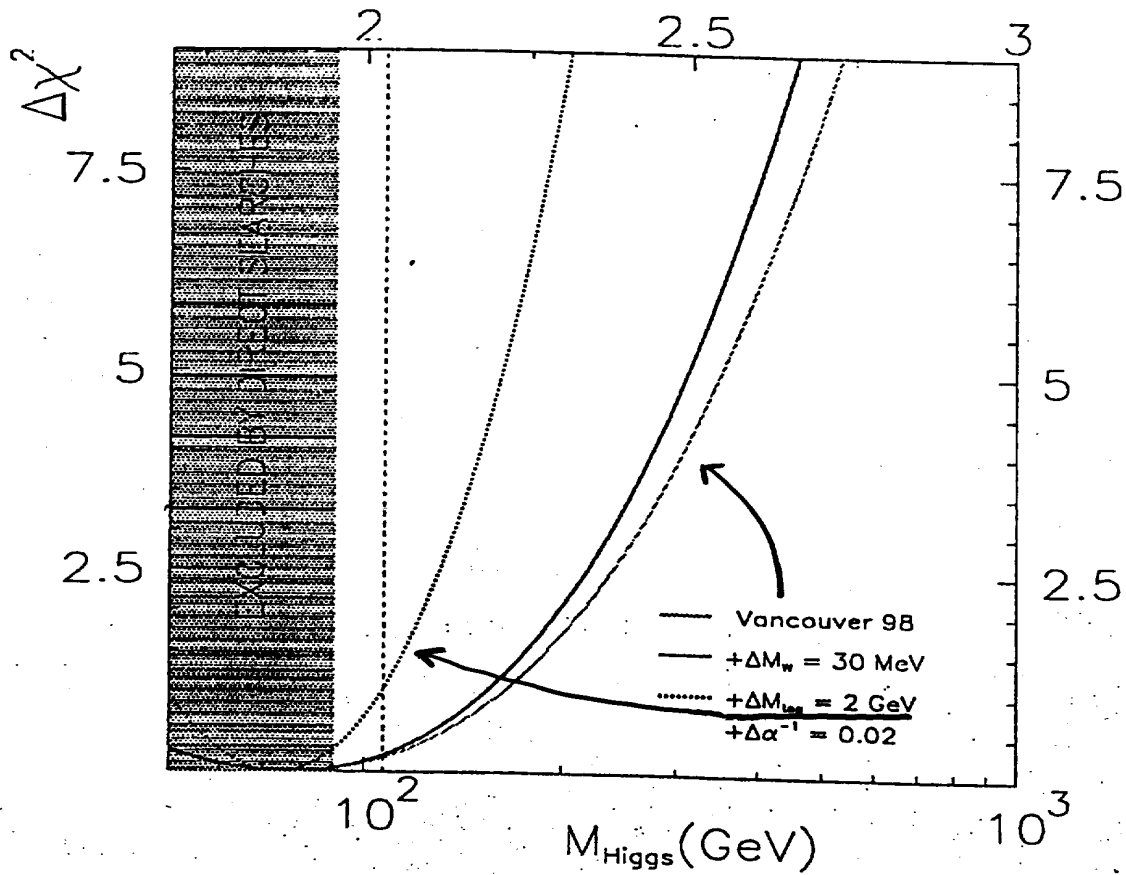
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$M_H < 262 \text{ GeV @ } 95\% \text{ C.L.}$

Teubert
RADCOR '98



Exp. accuracy [present and expected]

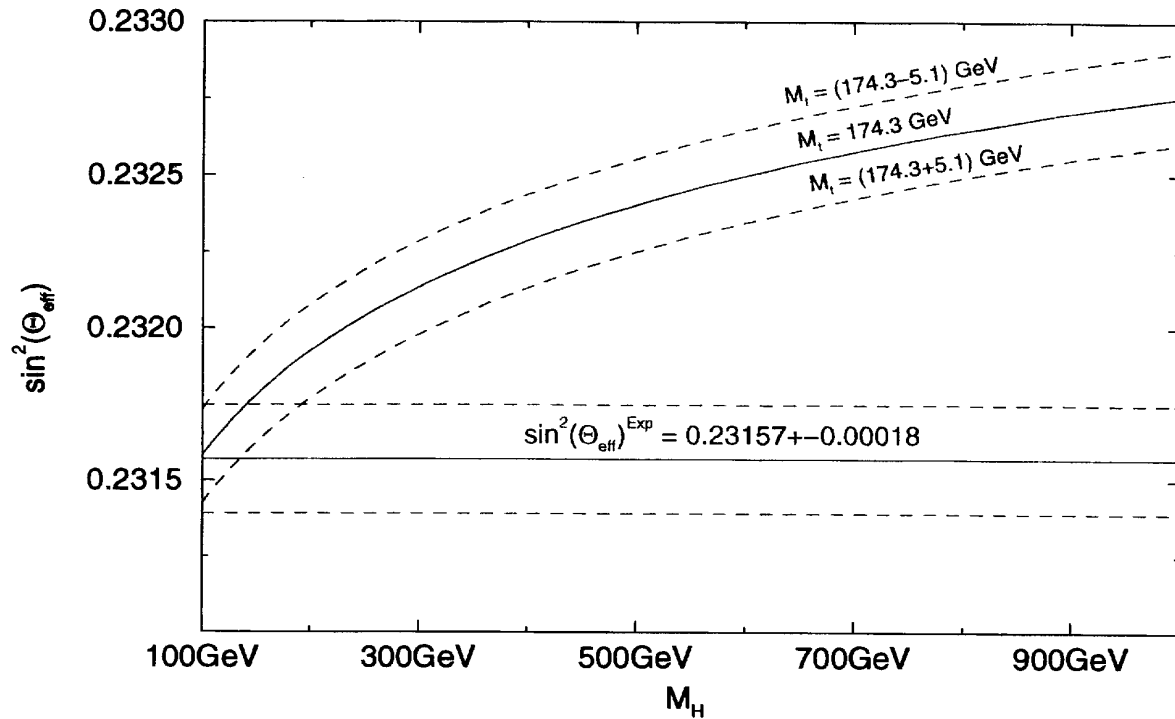
	LEP2 / TeV	LHC
M_W	40 MeV	15 - 20 MeV
$\sin^2 \theta_e$	0.00018	0.00018
m_t	5 GeV	2 GeV

work done with

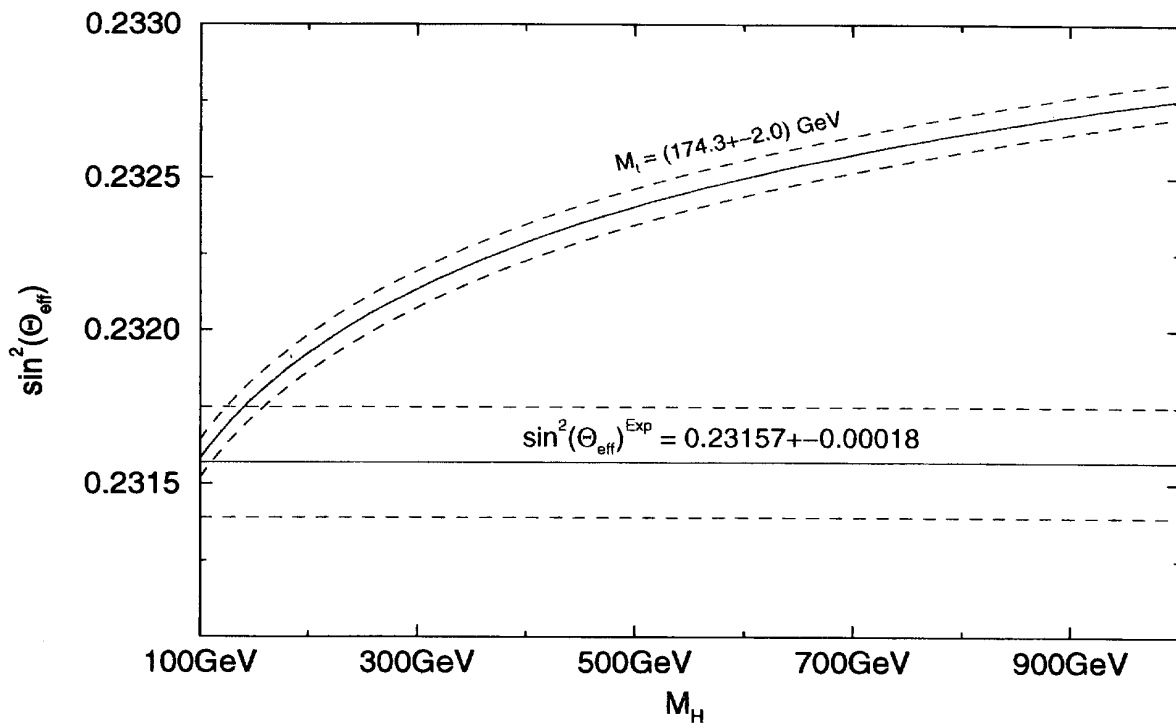
S. Heinemeyer, G. Weiglein

Higgs-mass dependence of $\sin^2 \theta_{\text{eff}}$ in the SM

Current experimental precision:

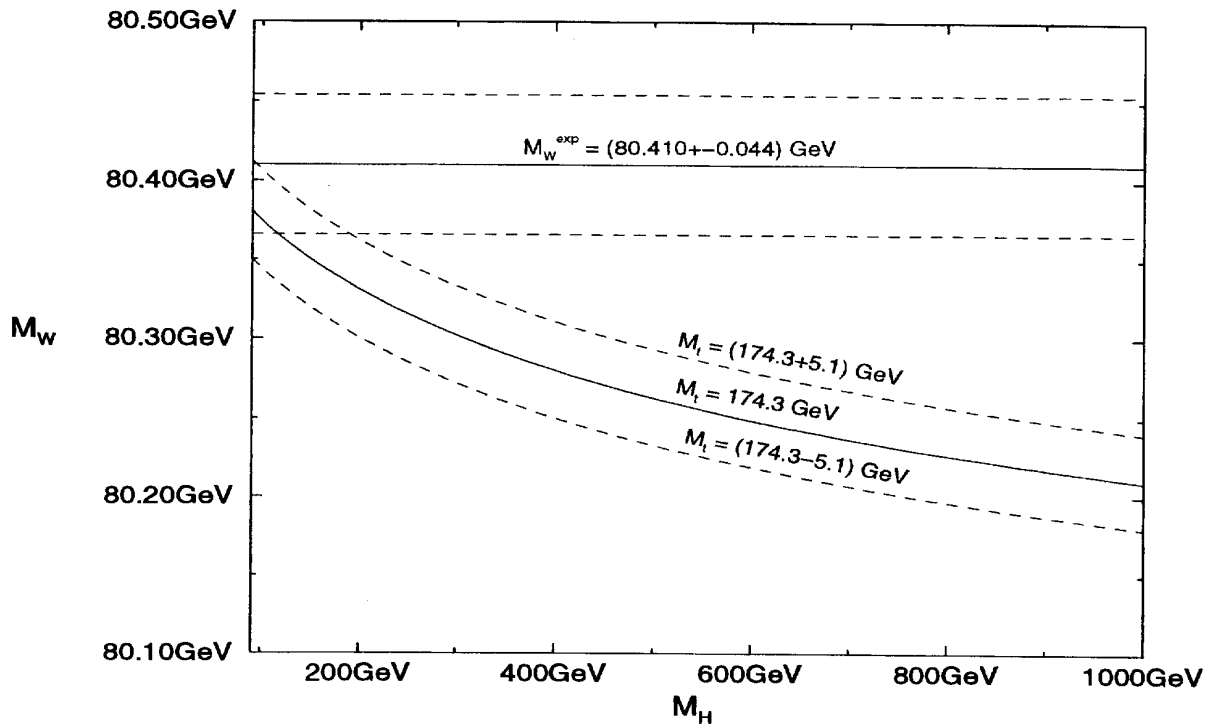


Experimental precision at LHC:

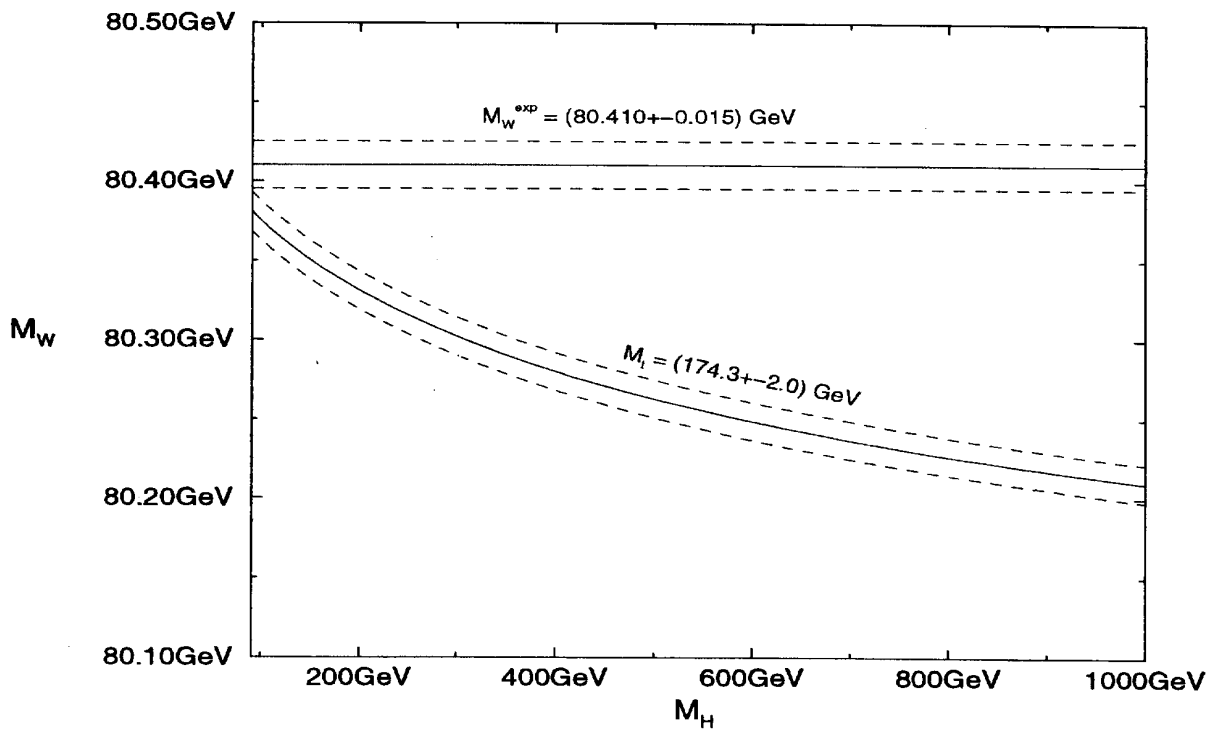


Higgs-mass dependence of M_W in the SM

Current experimental precision:



Experimental precision at LHC:



	M_H			$\delta m_t = 2 \text{ GeV}$	$\delta \alpha_h^*)$	δh_0
	120	150	180			
M_W	80.363	80.349	80.337	0.012	0.013 0.005 0.003	0.004
$\sin^2 \theta_e$	0.23167	0.23179	0.23188	0.00006	0.00023 0.00008 0.00005	0.00004

S_e	1.0051	1.0050	1.0048	0.0002	0.00005
R_b		0.2158		0.0001	

$$\delta S_e^{\text{exp}} = 0.0012$$

$$\delta R_b^{\text{exp}} = 0.00074$$

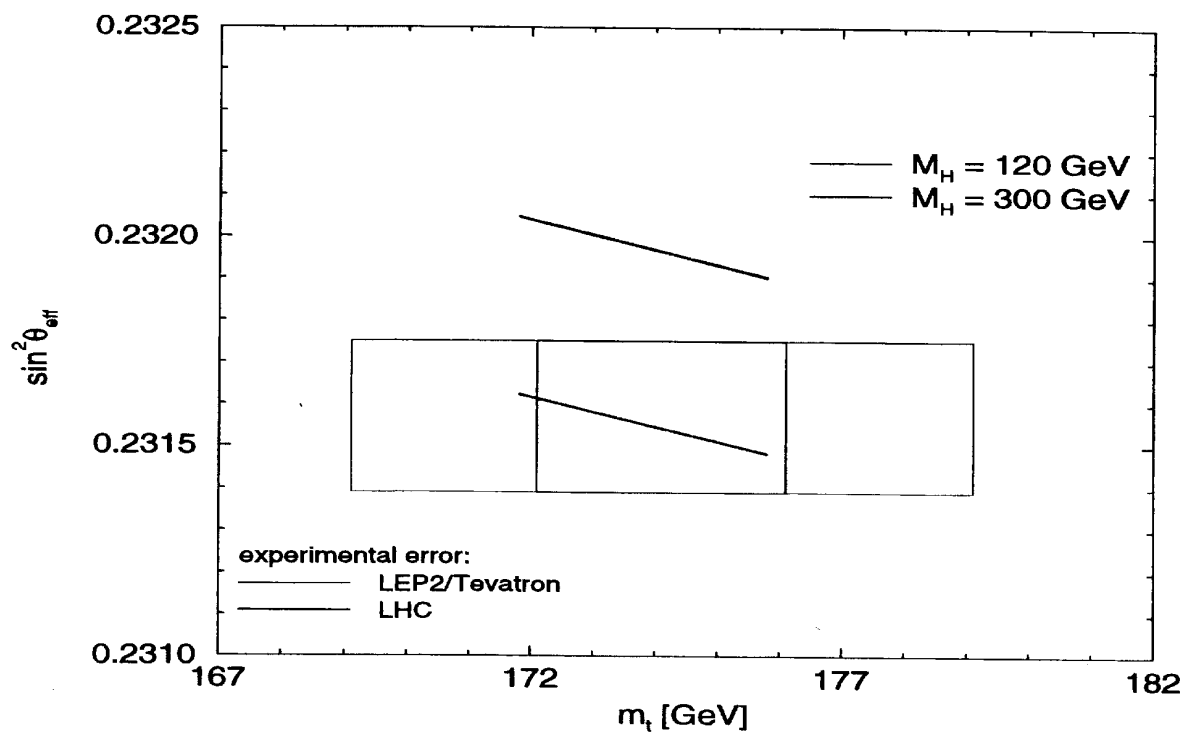
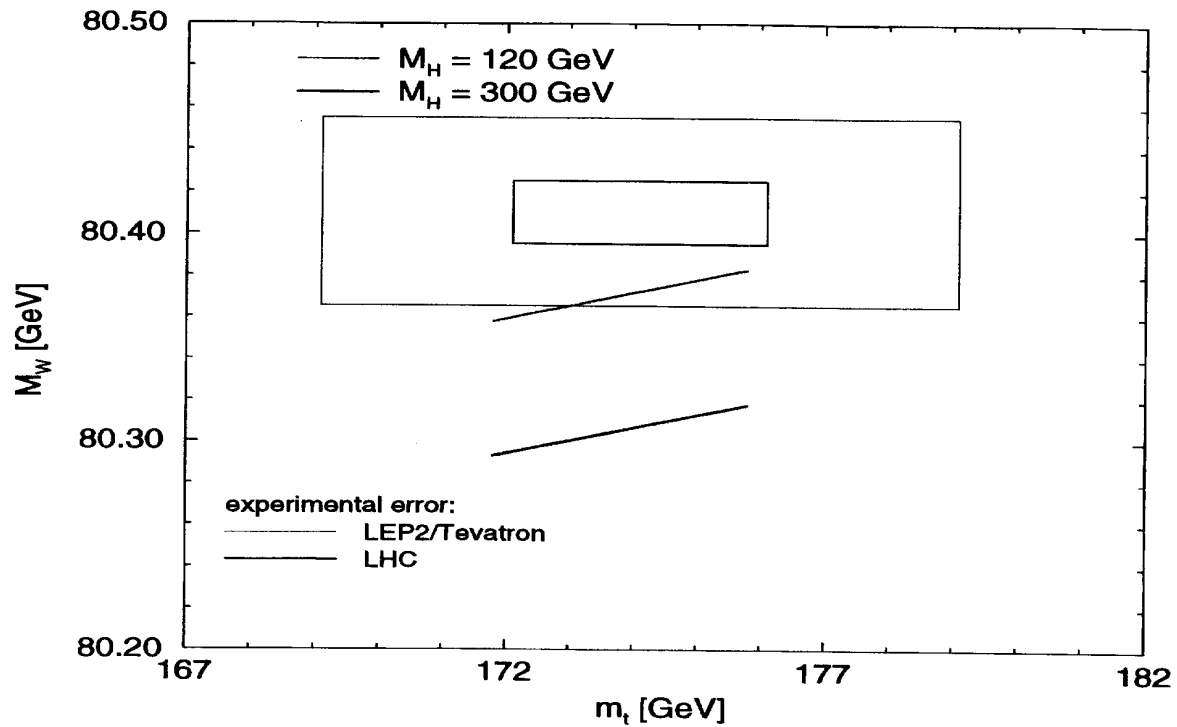
$$* \delta(\Delta\alpha) = 0.0007$$

$$= 0.000254$$

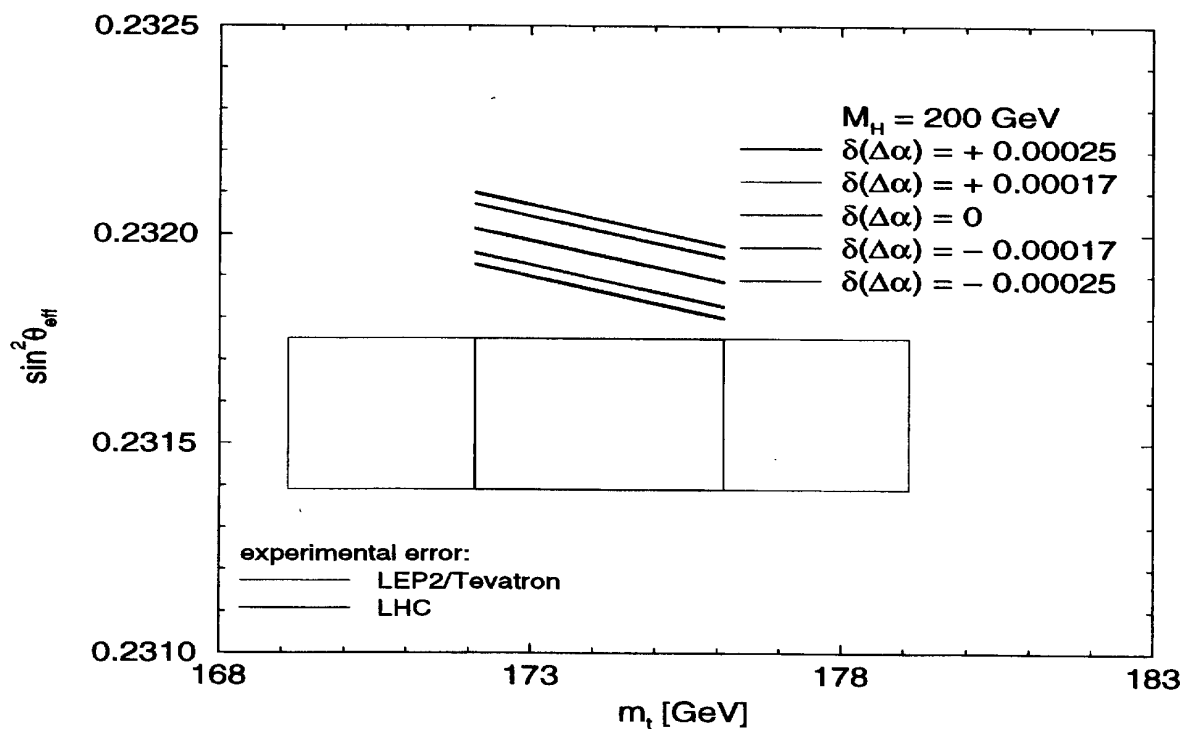
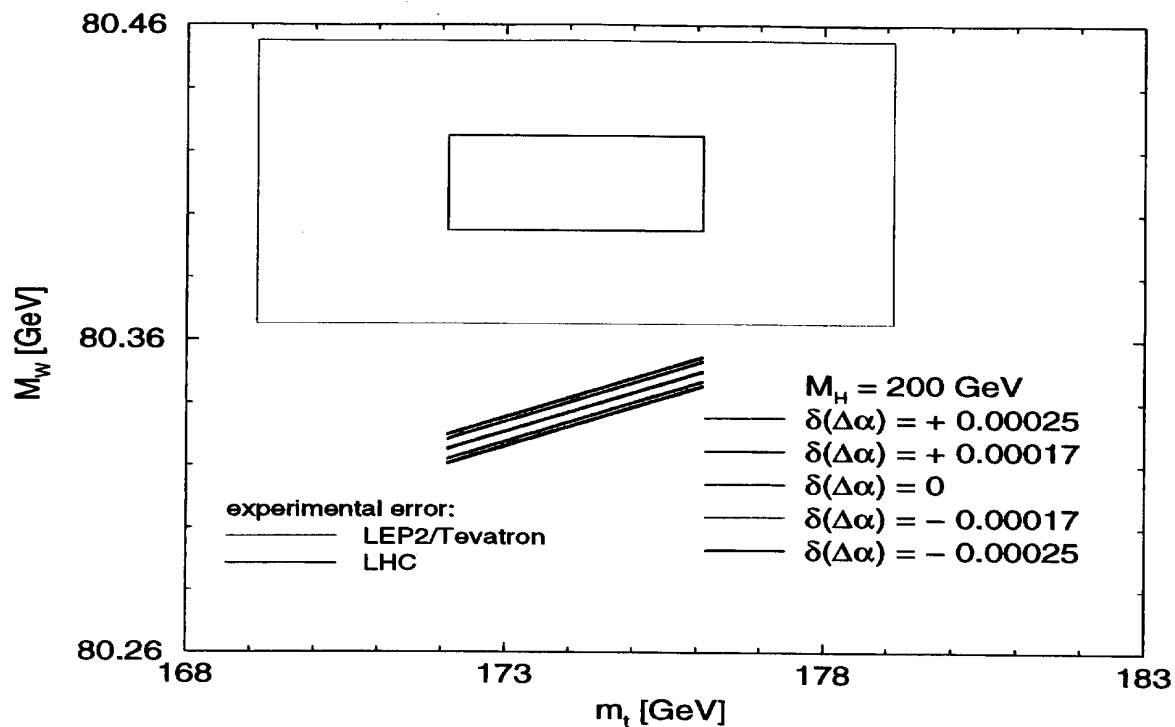
$$= 0.00017$$

Prediction for M_W , $\sin^2 \theta_{\text{eff}}$ in the SM,
 experimental precision: LEP2/Tev., LHC

Sensitivity to variation of m_h :



Prediction for M_W , $\sin^2 \theta_{\text{eff}}$ in the SM
 Effect of uncertainty in $\Delta\alpha$:



determination of M_H :

$$M_W = f_w (M_H, m_t, \dots)$$

$$\sin^2 \theta_e = f_s (M_H, m_t, \dots)$$

invert: $M_H = f_w^{-1} (M_W, m_t, \dots)$

$$M_H = f_s^{-1} (\sin^2 \theta_e, m_t, \dots)$$

	M_W	$\sin^2 \theta_e$
LEP 2 / Tevatron	62 %	39 %
LHC *)	32 %	38 %

*) $\delta m_t = 2 \text{ GeV}, \quad \delta M_W = 15 \text{ MeV}, \quad \delta \alpha_h = 0.00017$

non-standard physics (virtual)

- new physics contributions to ρ
 e.g. new generation(s) of fermions
 extra Higgses (doublet, triplet, ...)
 new dynamics (strongly int. sector)
 \Rightarrow weaker bound on M_H

- new physics to $\Delta\rho = \epsilon_1, \epsilon_2, \epsilon_3$ (T, U, S)
 bounds on M_H disappear

$$M_H = 390^{+690}_{-310} \text{ GeV} \quad \text{J. Erler (199)}$$

$$m_t = 172.9 \pm 4.8 \text{ GeV}$$

- Specific models: supersymmetry

MSSM \Rightarrow precision calculations
of EW observables

Test of extended models via extra contribution to $\Delta\rho$:

Sensitivity at LEP/Tevatron, LHC

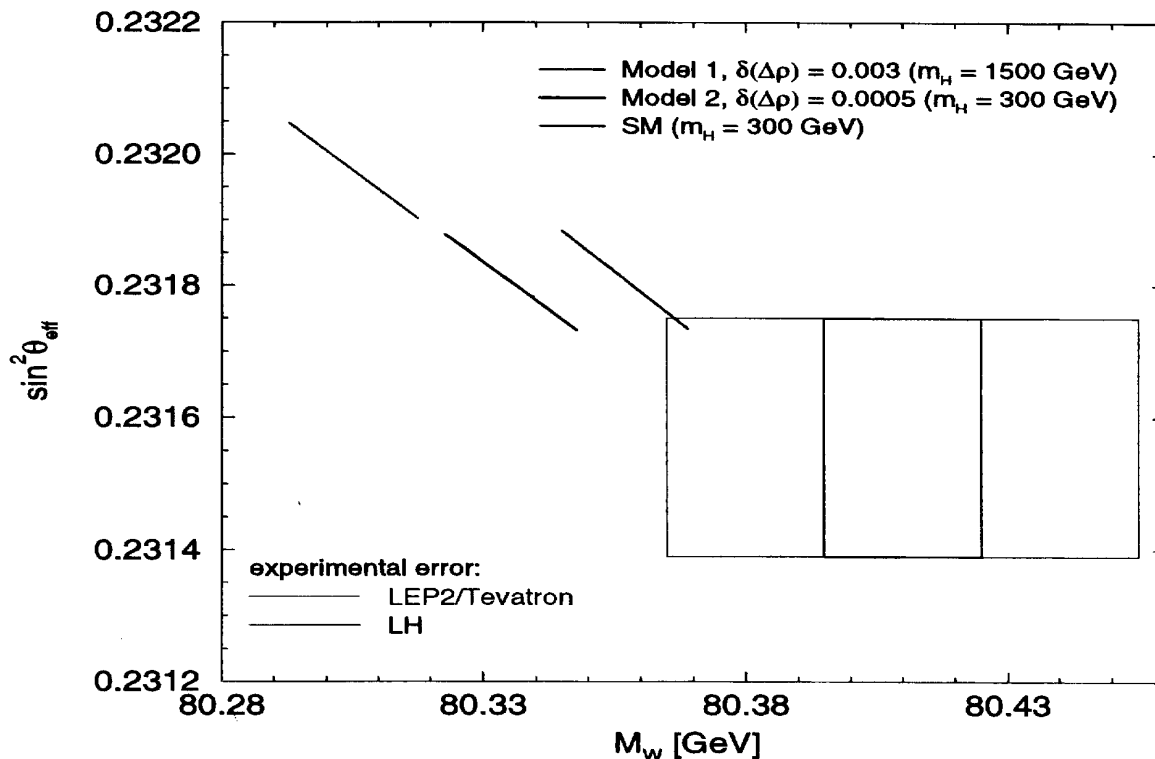
Two examples:

Model 1: Strong interacting Higgs bosons
 $\rightarrow M_H \gtrsim 900$ GeV (plot: $M_H = 1500$ GeV)
 $\Delta\rho^{extra} = \mathcal{O}(3 \times 10^{-3})$

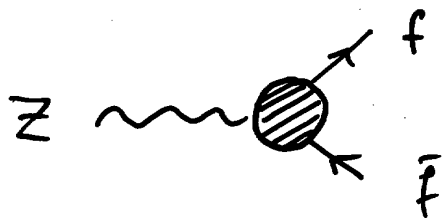
Model 2: Extra dimensions at low energies

[M. Masip, A. Pomarol '99]

$\Delta\rho^{extra} = \mathcal{O}(5 \times 10^{-4})$



- Effective Z boson couplings:



$$g_A^f = \sqrt{\rho_f} I_3^f$$

$$g_V^f = \sqrt{\rho_f} (I_3^f - 2 Q_f \sin^2 \theta_f)$$

ρ_f ($m_t, \alpha_s, M_A, \tan\beta$, SUSY-parameters)

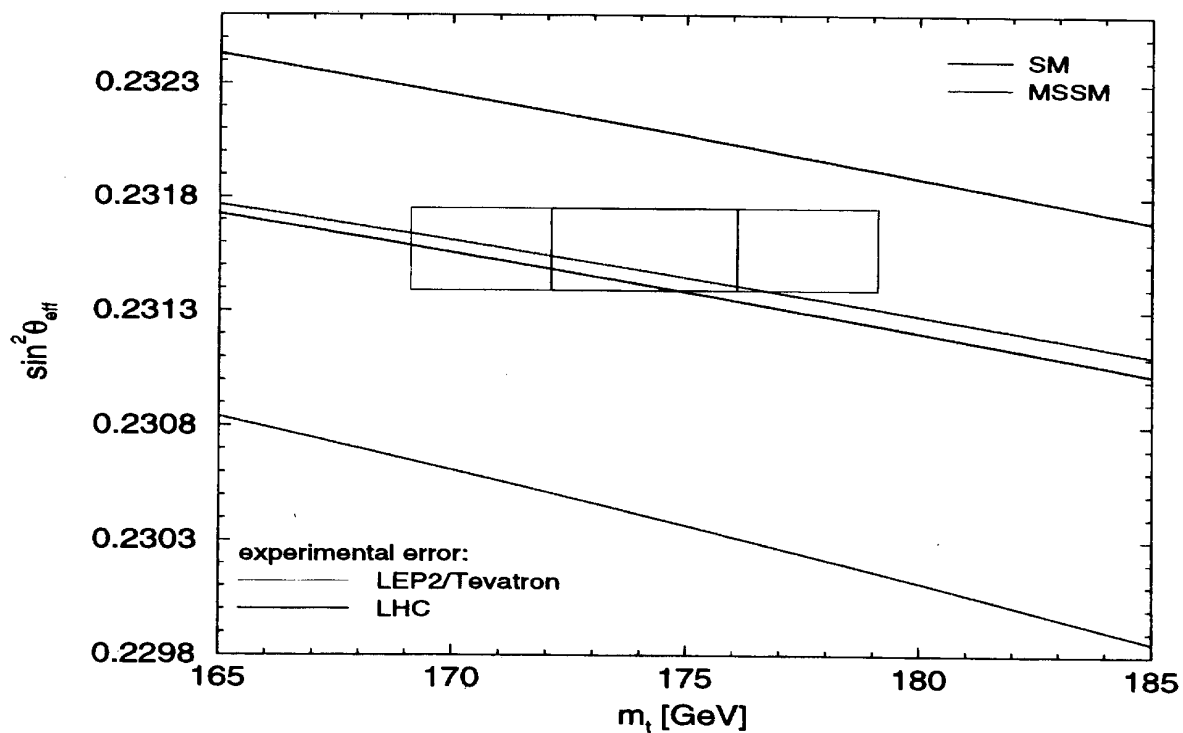
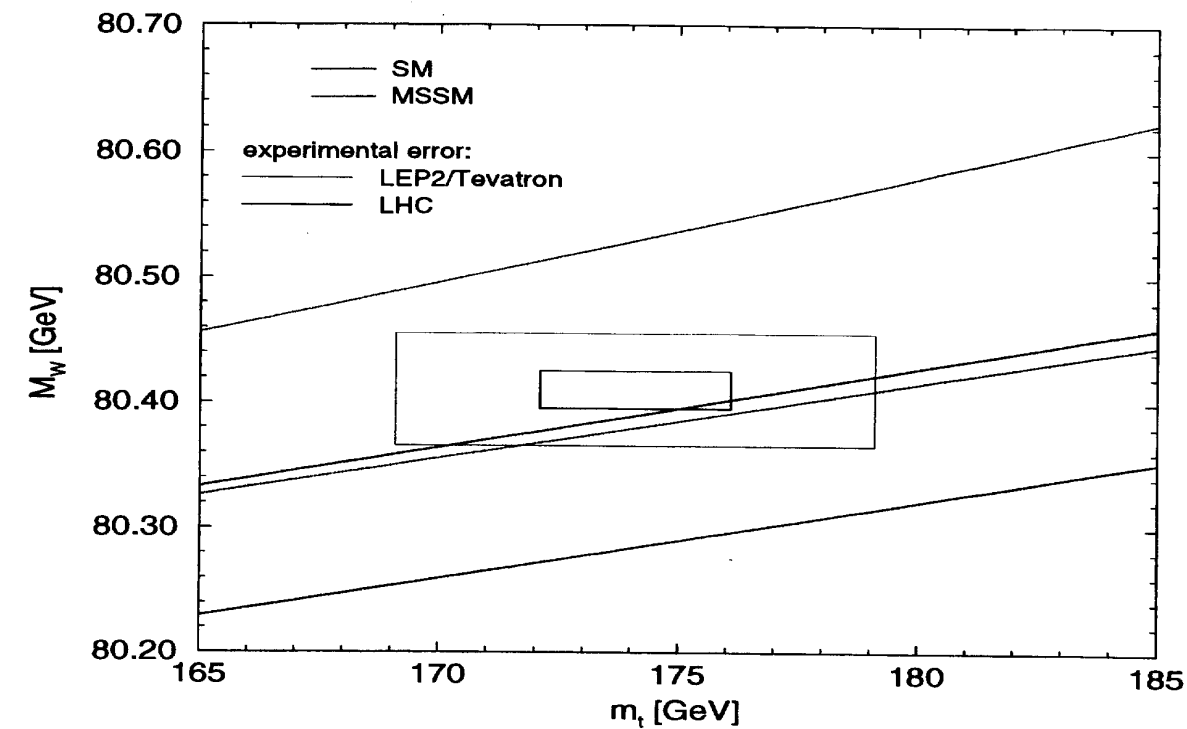
$\sin^2 \theta_f$ ($m_t, \alpha_s, M_A, \tan\beta$, SUSY-parameters)

- $M_W - M_Z$ correlation:

$$M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right) = \frac{\pi \alpha}{\sqrt{2} G_F} \cdot \frac{1}{1 - \Delta r}$$

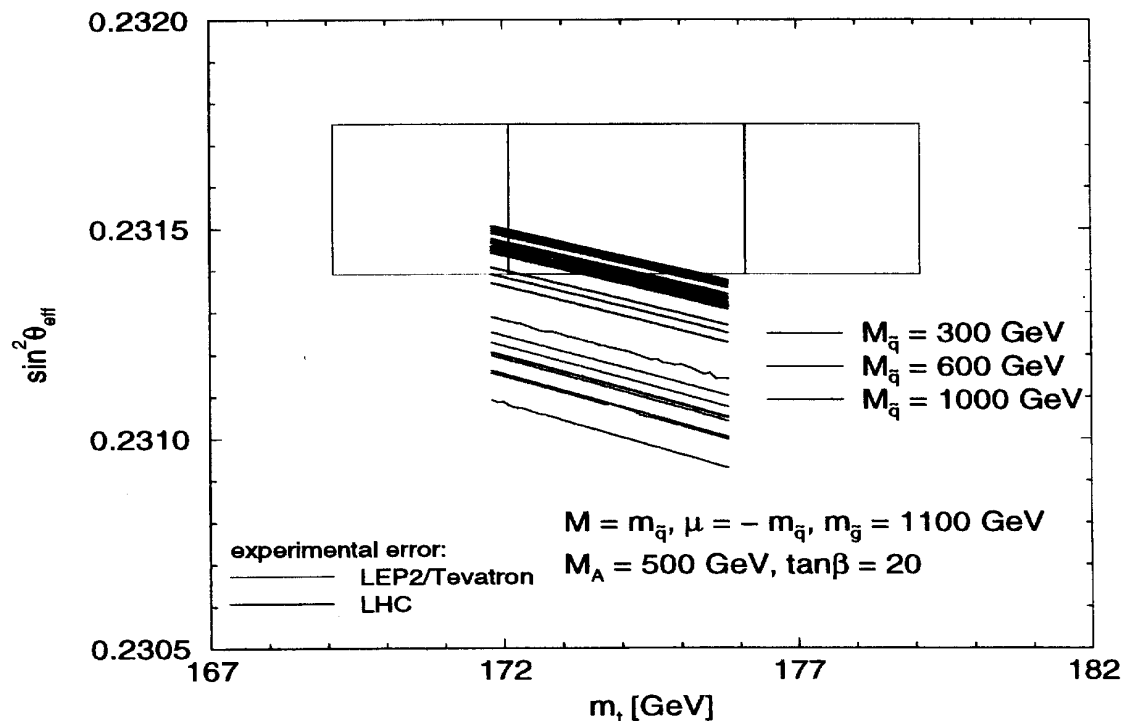
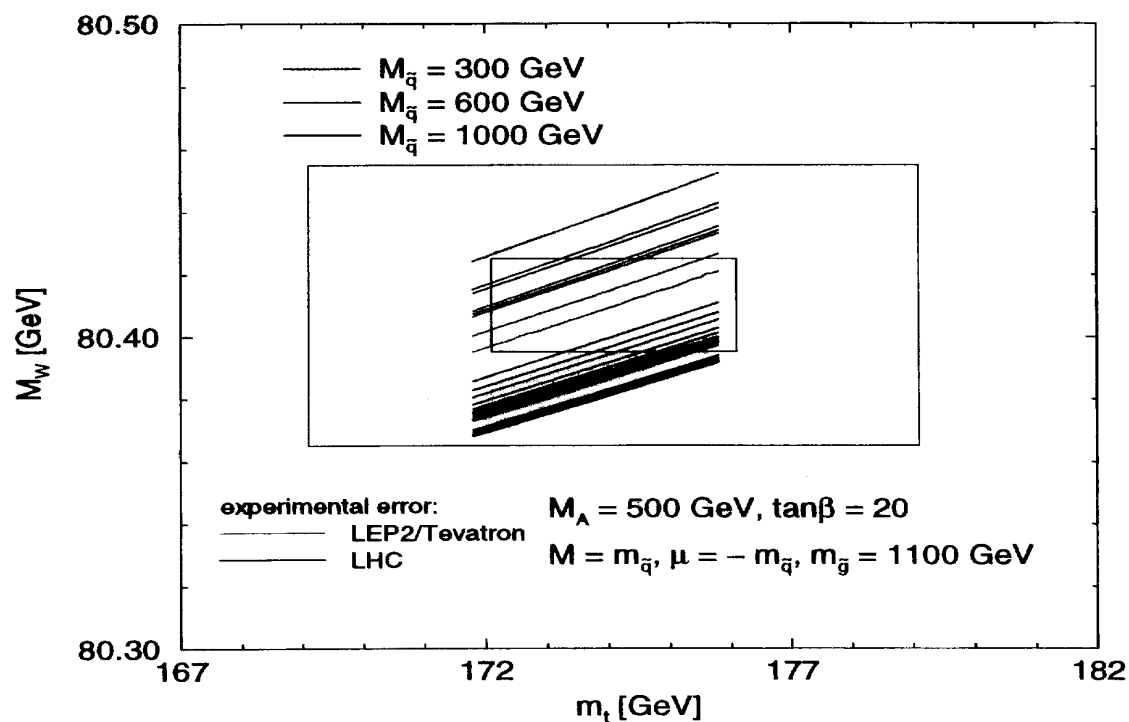
$\Delta r = \Delta r(m_t, \alpha_s, M_A, \tan\beta, \text{SUSY-par.})$

Prediction for M_W , $\sin^2 \theta_{\text{eff}}$ in SM and MSSM,
 experimental precision: LEP2/Tev., LHC



Prediction for M_W , $\sin^2 \theta_{\text{eff}}$ in the MSSM,
 experimental precision: LEP2/Tev., LHC

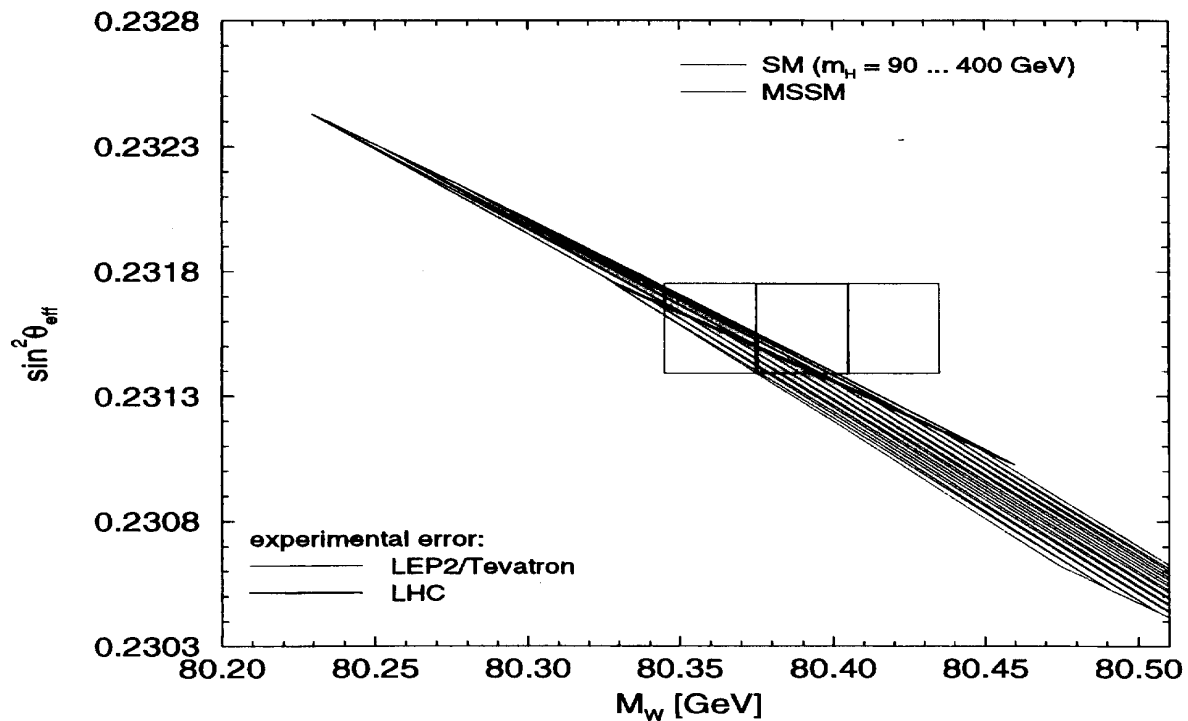
Sensitivity to variation of $m_{\tilde{q}}$ by $\pm 10\%$ and of
 mixing in \tilde{t} -sector:



Prediction for M_W , $\sin^2 \theta_{\text{eff}}$ in SM and MSSM,
 experimental precision: LEP2/Tev., LHC

SM: $90 \text{ GeV} \leq m_h \leq 400 \text{ GeV}$, m_t fixed 165–185

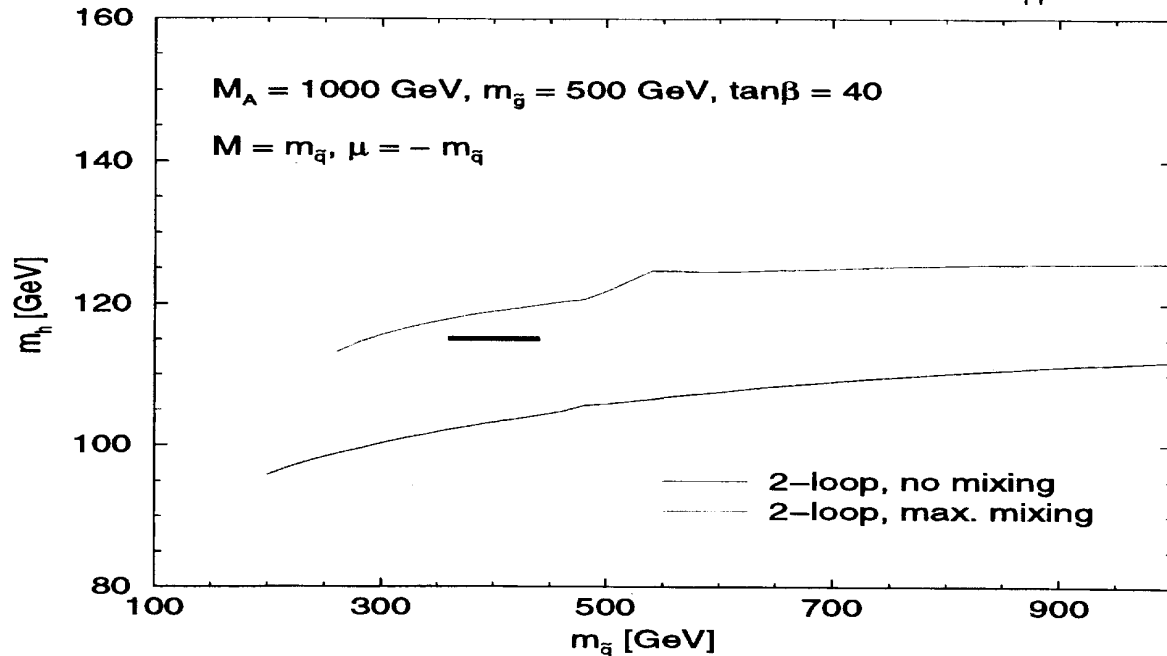
MSSM: $m_{\tilde{q}}$, mixing in \tilde{t} -sector varied



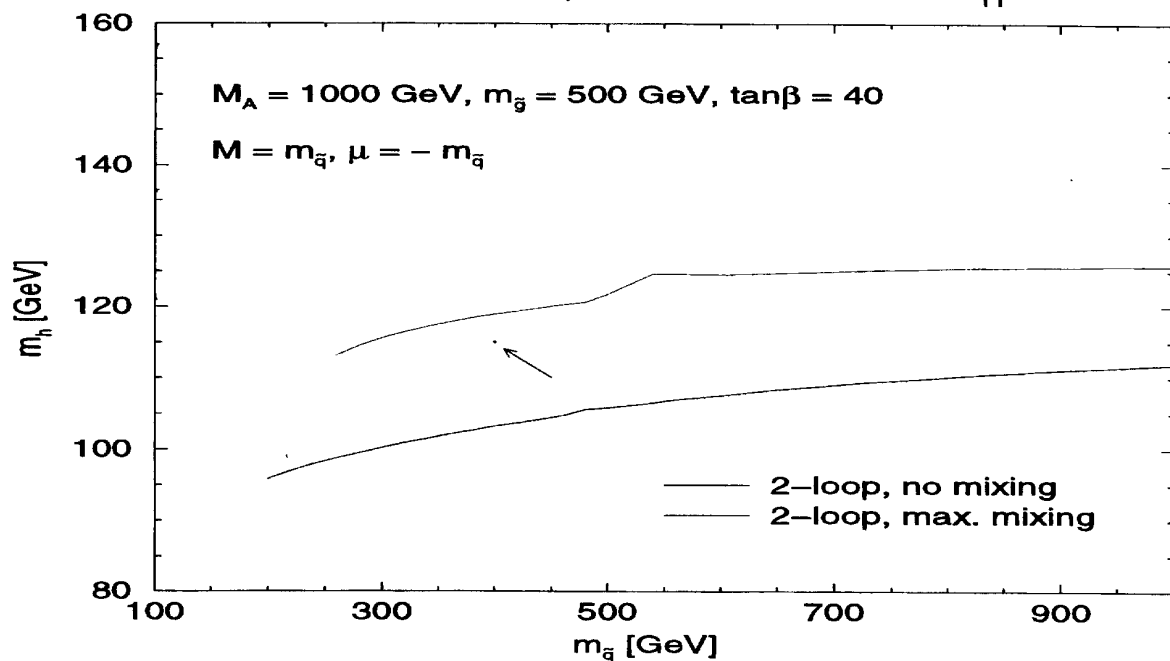
LHC/LC: Experimental information on m_h , M_{SUSY}

Assume: $m_h^{\text{exp}} = 115 \text{ GeV}$, $M_{\text{SUSY}}^{\text{exp}} = 400 \text{ GeV}$

LHC results vs. MSSM prediction for m_h :

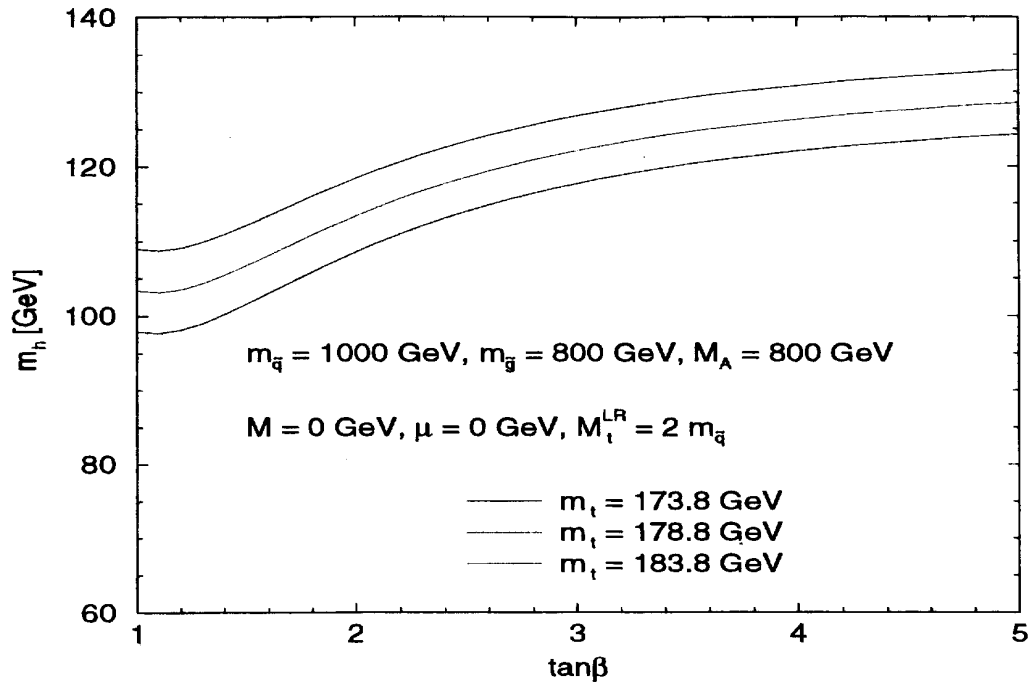


LC results vs. MSSM prediction for m_h :



⇒ Stringent test of MSSM

Maximal m_h -value in scenario with small $\tan \beta$,
 $m_{\tilde{q}} = 1000$ GeV for different values of m_t :
 $m_t = m_t^{\text{exp}}, m_t^{\text{exp}} + 1\sigma, m_t^{\text{exp}} + 2\sigma$



Values increase by 3–4 GeV for $m_{\tilde{q}} = 2000$ GeV

LEP2 exclusion limit: ≈ 107 GeV

$\Rightarrow m_h^{\text{max}}$ for $\tan \beta = 1.6$ is at the edge of
 LEP2 reach

$\tan \beta$ reach: $\tan \beta = 1.9$ for $m_t = 173.8$ GeV
 $\tan \beta = 1.5$ for $m_t = 178.8$ GeV
 for $m_{\tilde{q}} = 1000$ GeV