

# Triple Gauge-boson Couplings in Boson pair production

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## Outline:

- Introduction

Why }  
      } What }  
      } How

- Present status

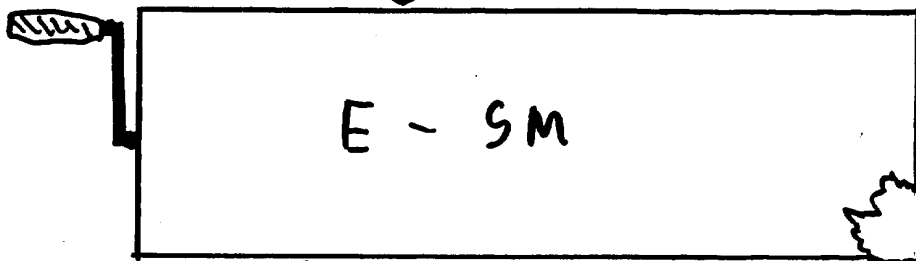
TeVatron

- At LHC / ATLAS

TGC

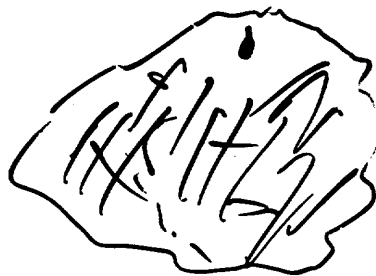
Fermion  
Zoology

Gauge Invariance  
 $SU_2 \times U_1$



Higgs  
chewing  
Gum

TGC  
• Gauge Group  
structure



Boson Zoology

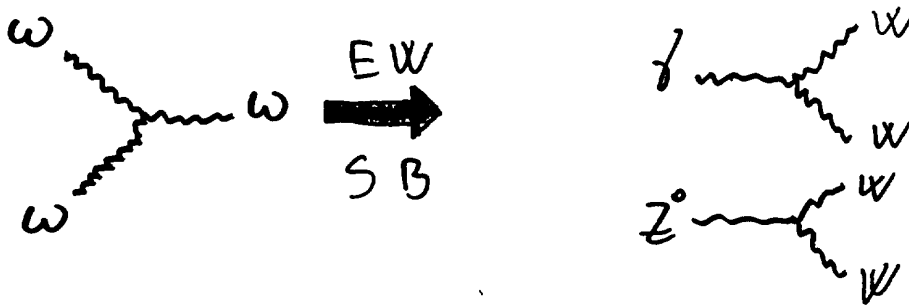


Boson properties  
Charge  
electric/magnetic  
moment  
Mass



Fermion interaction

The starting point for the TGCs is the non-Abelian structure of SM



Therefore:

- The TGCs are tightly connected with symmetry breaking
- The couplings of the  $\gamma W W$  and  $Z^0 W W$  vertices reflect the full gauge group structure
- Connected to  $W$  properties

$$Q_W = e g_1^Y \quad (\text{charge})$$

$$\mu_W = \frac{e}{2m_W} (g_1^Y + \kappa_Y + \lambda_Y) \quad (\text{Magn. dip.})$$

$$q_W = -\frac{e}{m_W^2} (\kappa_Y - \lambda_Y) \quad (\text{Elec. quad.})$$

➔ Test SM with  $\gamma W W$  and  $Z^0 W W$

- One step further (pessimistic scenario....)

TGCs also probe for New Physics

- In this context additional non-SM contributions has the magnitude

$$|TGC| \sim \left( \frac{M_U}{\Lambda_{NP}} \right)^2$$

- But: In order to have a completely model independent description we need 7 TGCs per vertex  $\rightarrow$  14 TGCs!  
Allows a determination of the structure of NP!
- Restricting us to C, P and em. gauge invariant TGCs:

$$\underbrace{g_1^Z, \kappa_\gamma, \kappa_Z}_{= 1 \text{ in SM}} \quad \text{and} \quad \underbrace{\lambda_\gamma, \lambda_Z}_{= 0 \text{ in SM}}$$

- What can we expect?

$$\Lambda_{NP} = 1 \text{ TeV} \Rightarrow TGCs \sim \mathcal{O}(10^{-2})$$

- Radiative corrections:  $\mathcal{O}(10^{-3})$
- SUGRA - Gut MSSM:  $\mathcal{O}(10^{-3})$
- Unconstrained MSSM:  $\mathcal{O}(10^{-2})$

## Characteristics of TGCs

- Main sensitivity lie in the angular distribution of the boson production angles  $\Rightarrow$
- But: different TGCs contribute to different boson helicity states
  - $\rightarrow$  Decay angular information enhance sensitivity to individual TGCs
  - $\rightarrow$  Advantage for clean  $e^+e^-$  machines
- Since the Lagrangian is linear in the TGCs we have

$$\frac{d\sigma}{d\Omega} = \underbrace{a_0}_{SM} + a_1 \cdot TGC + a_2 \cdot (TGC)^2$$

at most quadratic dependence!

- Generally we have, that

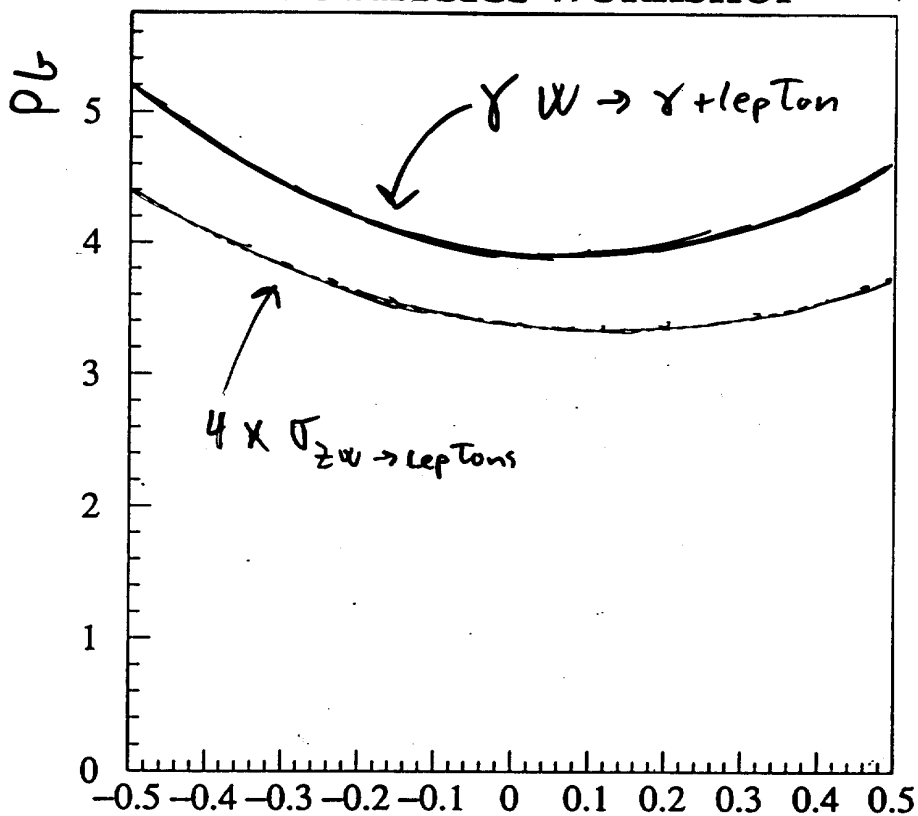
$$\sigma \propto + (TGC)^2 \quad (\text{Increasing})$$

- Sensitivity increases strongly with energy

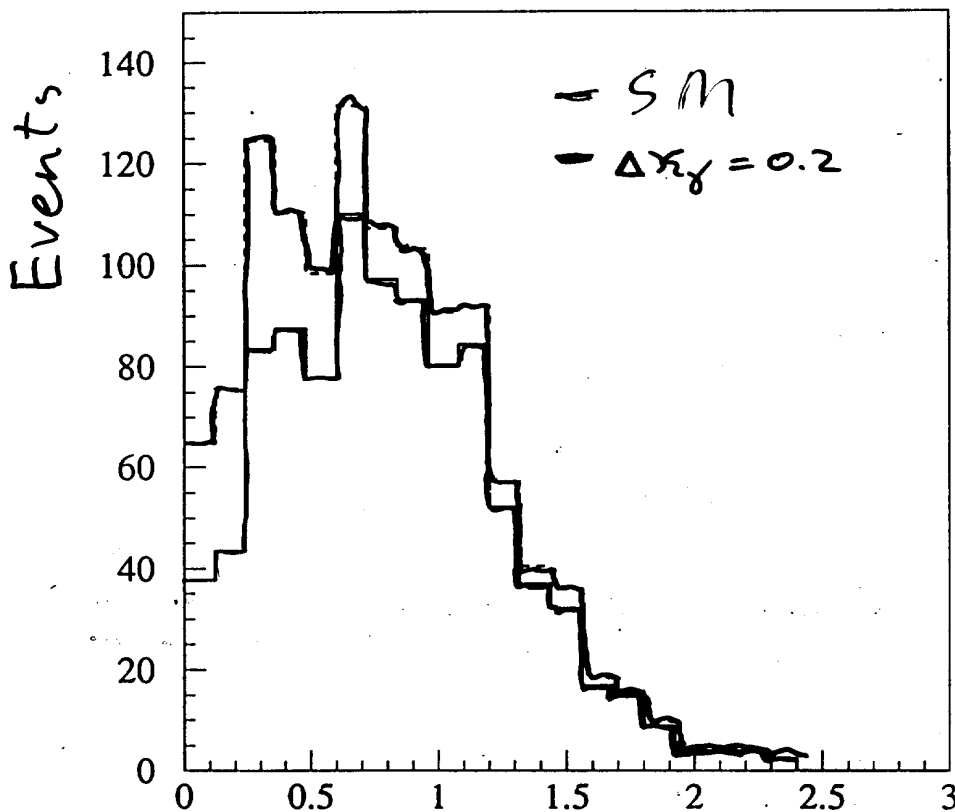
$$\frac{a_1 \cdot TGC + a_2 \cdot (TGC)^2}{a_0} \propto S$$

$\rightarrow$  Favor Hadron machines with their high energy ("p<sub>T</sub> counting")

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$\Delta \kappa_2$   
 $\Delta \kappa_2$



$$|y| = \frac{1}{2} \left| \ln \frac{1 + \cos \theta}{1 - \cos \theta} \right|$$

$\theta$ : angle between  $\gamma$  and beam in  $W\gamma$ -system

$|y|$

# TGCs from Tevatron

- At hadron colliders we have distinguishable final states

$$W\gamma, WZ$$

along with  $WW$

- Provide independent constraints on  $\gamma WW$  and  $ZWW$  couplings
- Important difference for hadron machines: The TGCs are introduced as form factors with scale  $\Lambda$

$$TGC \rightarrow \frac{TGC}{(1 + s/\Lambda^2)^2}$$

→ Higher  $\Lambda$  improve limits. Usually  $\Lambda \sim \sqrt{s}$

- At CDF/DØ all 3 final states are studied

- $W\gamma \rightarrow (e)\nu\gamma$

TGCs extracted via binned likelihood fit to  $E_T$  of  $\gamma$ .

- $WW \rightarrow l\nu l\nu$

Binned likelihood fit to  $E_T(l)$  and  $P_T(l\nu)$

- $WW/WZ \rightarrow l\nu ss/ll\gamma\gamma$  (high backg.)

Binned likelihood fit to  $P_T$  lepton system

$D\phi$  : Fermilab-Pub-97/136-E

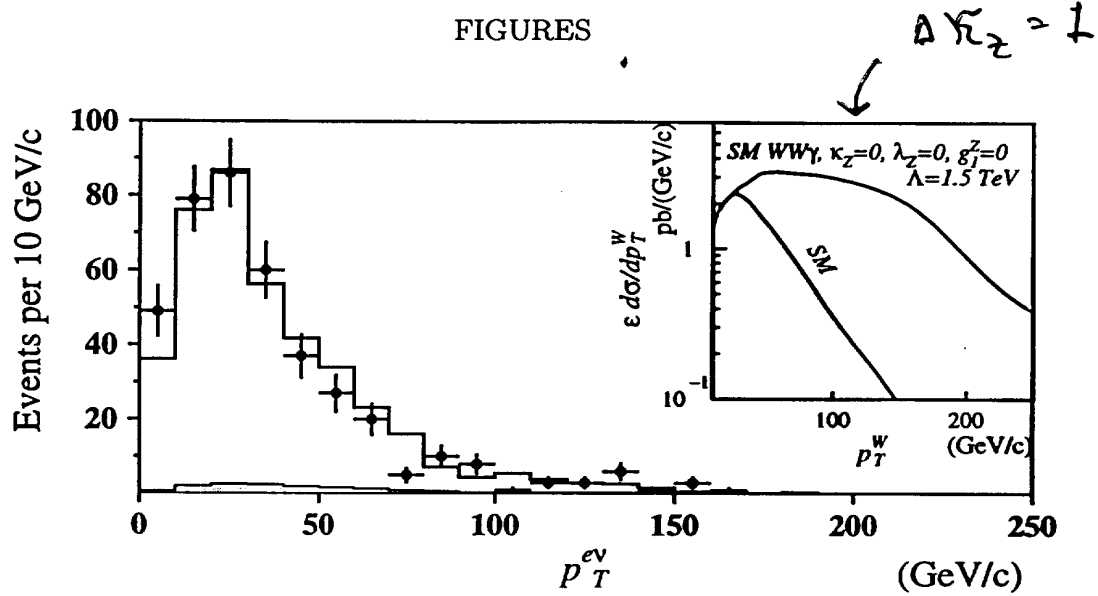
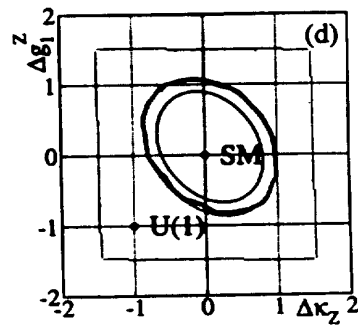
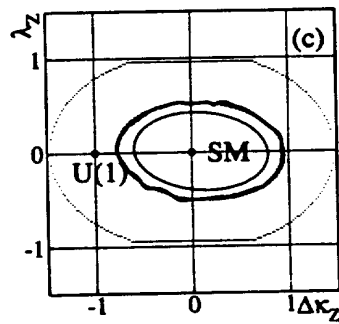
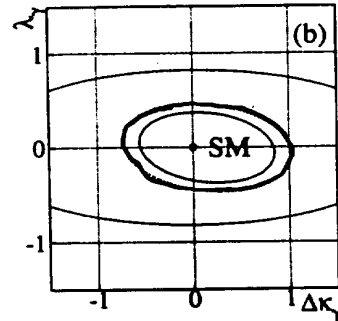
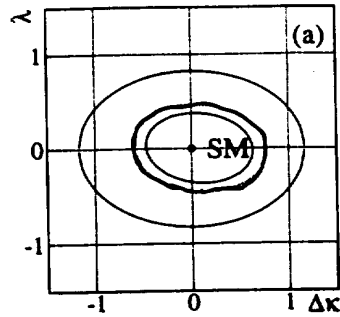


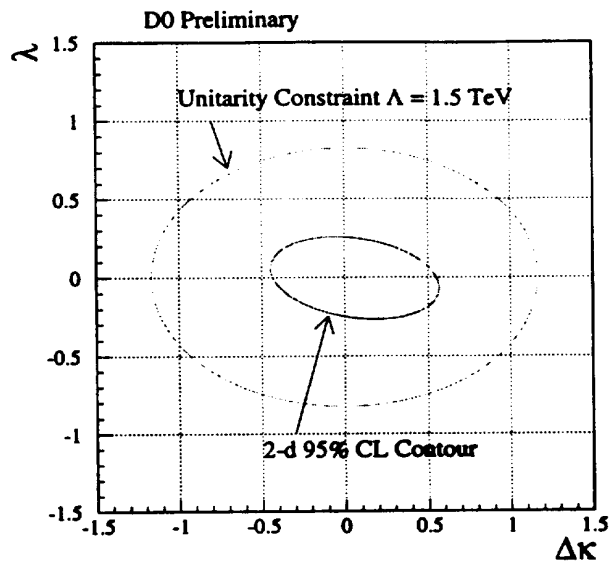
FIG. 1.  $p_T$  distributions of the  $e\nu$  system for the 1993–1995 data set. The points with error bars represent the data. The solid histogram is the total background estimate plus the SM Monte Carlo predictions of  $WW$  and  $WZ$  production (shown as shaded histogram). The inset shows the predicted  $d\sigma/dp_T^W$ , folded with the detection efficiencies, for SM  $WW\gamma$  and  $WWZ$  couplings (lower curve), and for SM  $WW\gamma$  and the indicated anomalous  $WWZ$  couplings (upper curve).



$D\phi$



- 95%



- Combined fits to all analyses gives

$$\begin{aligned}
 -0.93 < \Delta\kappa_\gamma < 0.94 \\
 -0.31 < \lambda_\gamma < 0.29 \\
 -0.33 < \Delta\kappa = (\Delta\kappa_1 = \Delta\kappa_2) < 0.45 \\
 -0.20 < \lambda = (\lambda_1 = \lambda_2) < 0.20
 \end{aligned}
 \quad
 \begin{aligned}
 & 95\% \text{ CL} \\
 & \Lambda = 2 \text{ TeV} \\
 & \text{run 1}
 \end{aligned}$$

To conclude

After so many years with SM we can still not constrain the Gauge group structure !

- Immediate future perspectives :

Ultimate limits from

LEP :  $\sim 3$  lower

Tevatron :  $\sim 3$  lower

→ 4-5 times lower

→  $\mathcal{O}(0.07)$  when LHC starts

- Scaling to  $10 \text{ fb}^{-1}$  @ 14 TeV should give  $\sim 10$  times lower !

## An ATLAS study of the Potential limits @ LHC

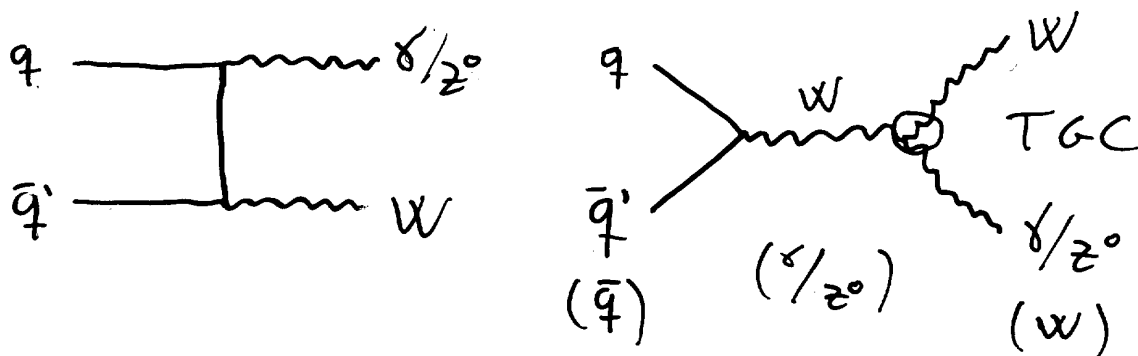
- Follows the line from the previous study (TP: 1994):
  - Leptonic final states for  $\gamma W$  and  $W Z$
  - Limits derived from high  $P_T$  event counting
  - AT  $\sqrt{s} = 14 \text{ TeV}$  and  $\int \mathcal{L} dt = 100 \text{ fb}^{-1}$  :

$$\begin{array}{l} \gamma W \\ \hline |\Delta\kappa_\gamma| < 0.07 \\ |\lambda_\gamma| < 0.0025 \end{array}$$

$$\Lambda = 10 \text{ TeV}$$

$$\begin{array}{l} Z W \\ \hline |\Delta\kappa_Z| < 0.13 \\ |\lambda_Z| < 0.006 \end{array}$$

- In hadron machines the dominant Gauge boson pair production processes are



- Considering lepton final states only  
the dominant physics background is  
(as usual)

$$pp \rightarrow t\bar{t} \rightarrow \text{lepton}(s) + X$$

$$pp \rightarrow b\bar{b} \rightarrow \text{lepton}(s) + X$$

$$pp \rightarrow \gamma b \rightarrow \gamma + \text{lepton} + X \quad (\gamma W)$$

$$pp \rightarrow \gamma z \rightarrow \gamma + \text{lepton}(s) + X \quad (\gamma W)$$

(not included yet:  $zQ\bar{Q}$ ,  $zQ$  + fake id  
combinatorics)

- Selection based on Isolated lepton /  
photon id + event quantities

$\gamma W$	$z W$
<p><u>One</u> lepton with <math>E_T &gt; 40 \text{ GeV}</math></p> <p>No <math>l\bar{l}</math> pair with <math>-15 &lt; M_{l\bar{l}} - M_z &lt; 15</math></p> <p><math>E_T^\gamma &gt; 100 \text{ GeV}</math></p> <p><math>M_T^{lv} &gt; 35 \text{ GeV}</math></p> <p>Extra leptons: <math>E_T &lt; 10 \text{ GeV}</math></p> <p>Jet energy veto <math>E_T^{\text{2nd jet}} &lt; 20 \text{ GeV}</math></p>	<p><u>Three</u> leptons with <math>E_T &gt; 10 \text{ GeV}</math> and <u>at least</u> one valid <math>l\bar{l}</math> flavour comb. <u>At least</u> one <math>l\bar{l}</math> pair with</p> <p><math>-20 &lt; M_{l\bar{l}} - M_z &lt; 20</math></p> <p><math>\text{Max}(E_T^l, E_T^{\bar{l}}) &gt; 25 \text{ GeV}</math></p> <p><math>\text{min}(E_T^l, E_T^{\bar{l}}) &gt; 15 \text{ GeV}</math></p> <p><math>E_T^{lv} &gt; 10 \text{ GeV}</math></p> <p><math>M_T^{lv} &gt; 35 \text{ GeV}</math></p>

- More sophisticated selection foreseen (eg. anti-b tag) but present selection is compatible with previous study and Tevatron
- Resulting event sample

$\gamma W$  :  $P = 94\%$  (caution)

Number of events/year (low)  $\sim 1000$

$Z W$  :  $\epsilon = 34\%$   $P = 84\%$  (caution)

Number of events/year (low)  $\sim 1500$

### Extracting TGCs

- The TGCs are extracted individually for  $\gamma W$  and  $Z W$  using unbinned maximum likelihood fit to the distribution of

$\gamma W$  :  $P_T^\gamma$

$Z W$  :  $P_T^Z$

- The likelihood for each event is evaluated from a MC distribution generated with Baur + Zeppenfeld  $\gamma W$  and  $Z W$  processes, which have been implemented into Pythia/ALPheg

- The advantage of this method is
  - To the extent that MC will describe reality -
  - that it automatically include detector and background effects.



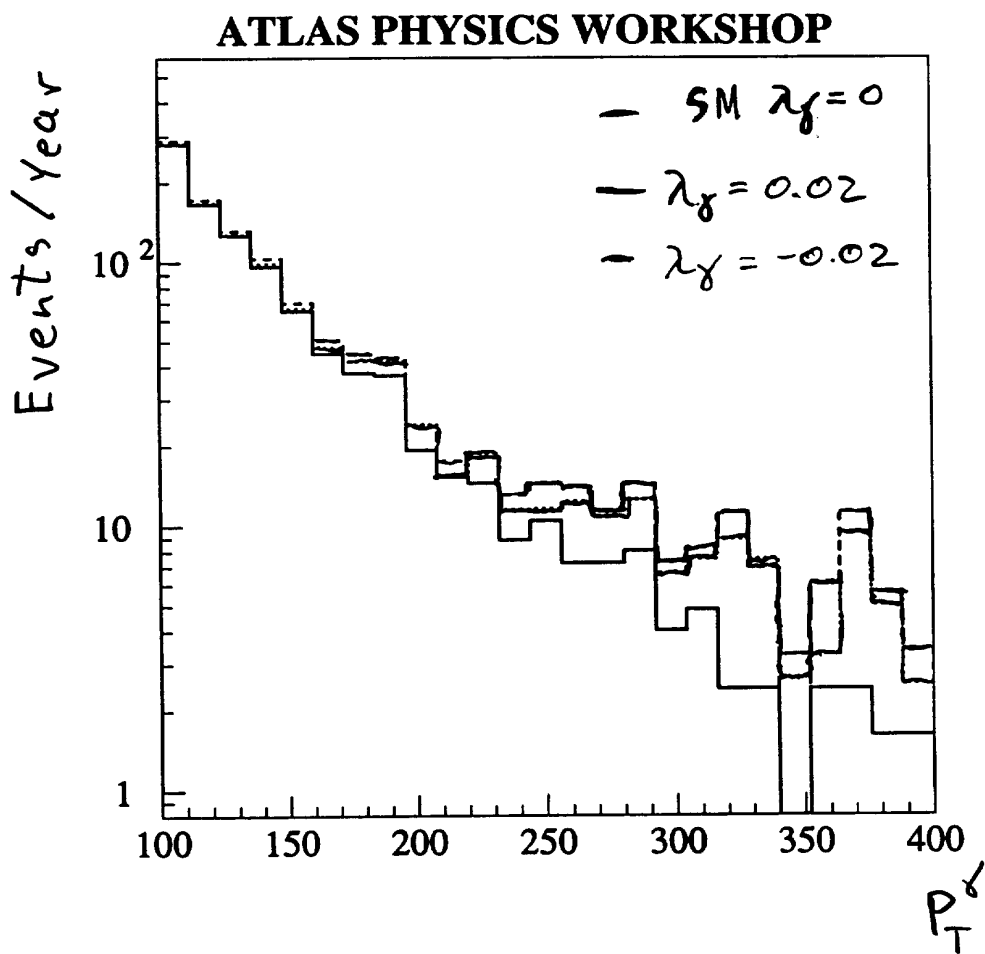
Equivalent to Tevatron

- Results ( $\sqrt{s} = 14 \text{ TeV}$ )

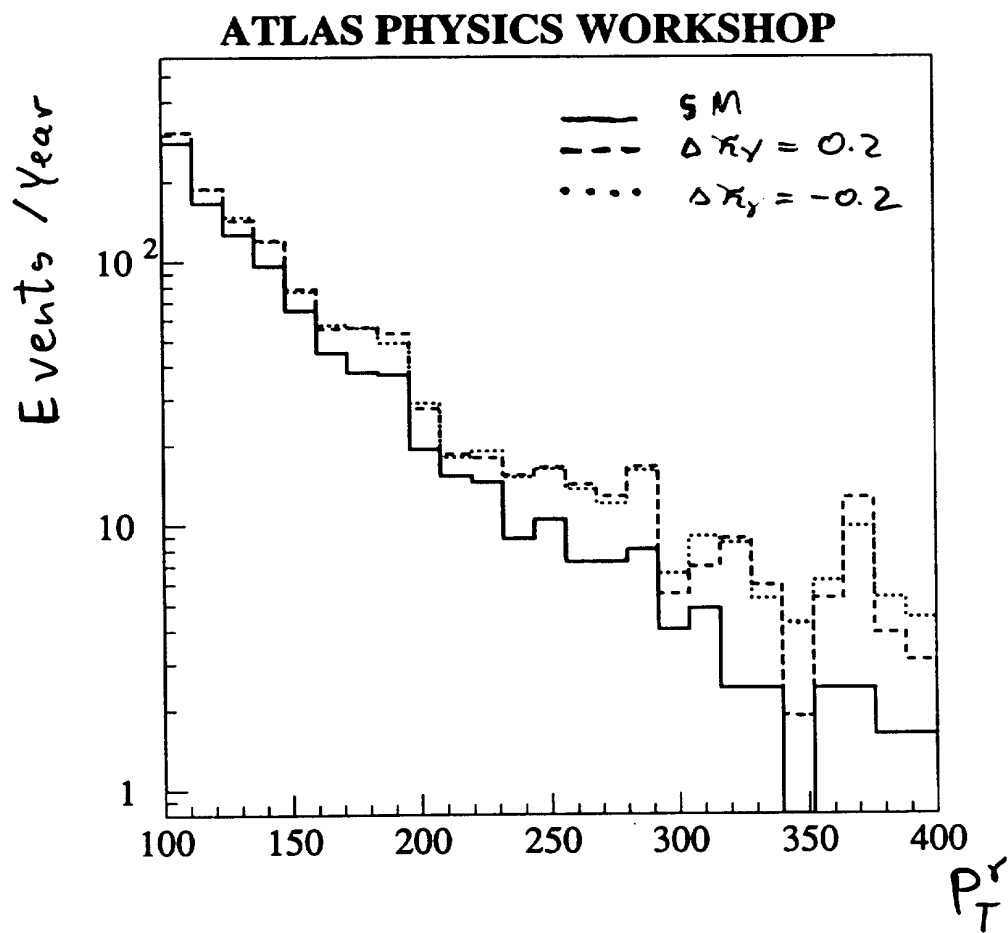
$0.063 < \Delta\kappa_\gamma < 0.060$
$0.005 < \lambda_\gamma < 0.007$
$0.087 < \Delta\kappa_2 < 0.129$
$0.013 < \lambda_2 < 0.028$

95% CL

- NOTE: Snapshot of ongoing studies!
  - Plan: Combine  $\gamma W$  and  $ZW$
  - Include angular information
  - Look into  $WW$  channel
  - High statistics background study

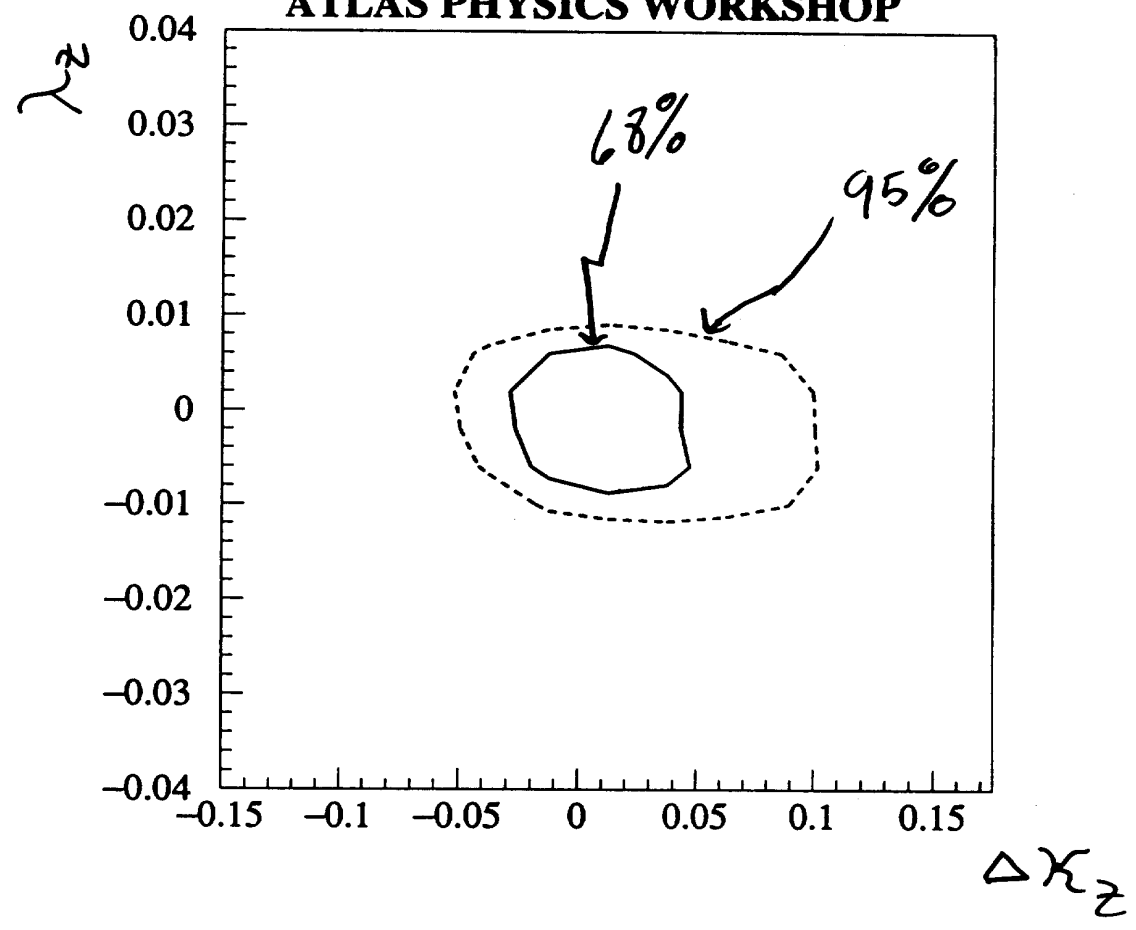
$\delta W$ 

- Similarity of opposite sign come from quadratic behaviour of total cross-section

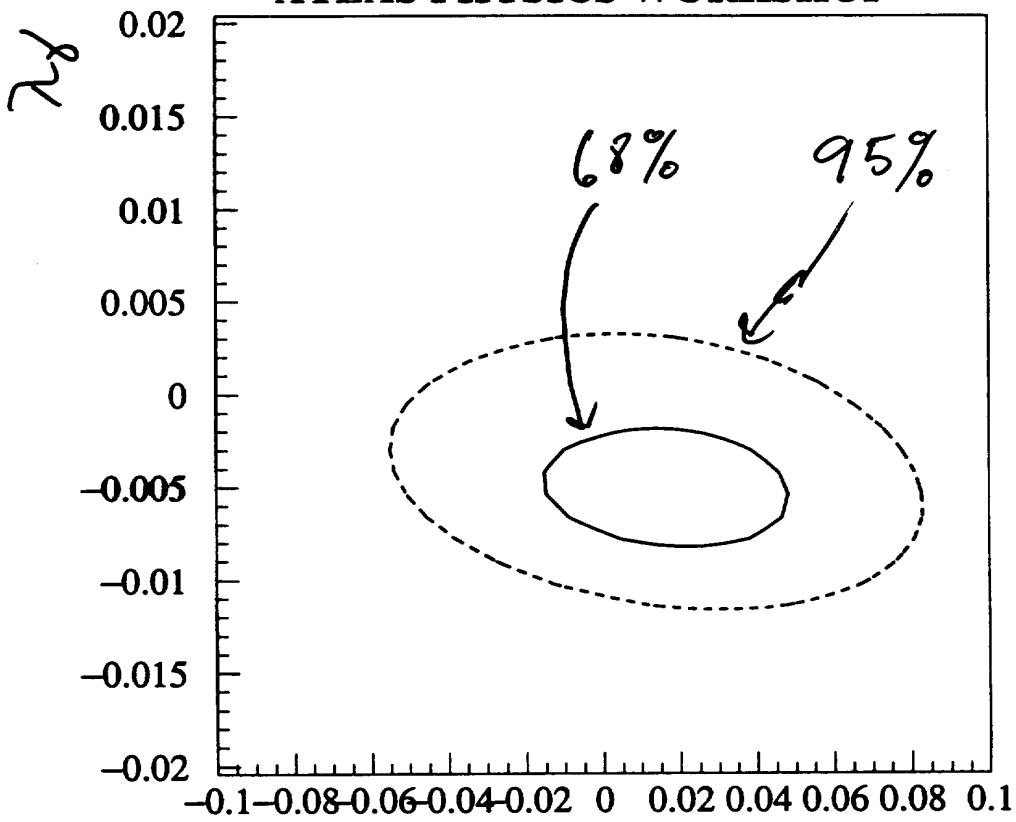
$\gamma W$ 



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$\Delta k$

## Conclusion

- The study of boson pair production provides a Unique determination of the SM Gauge group structure
- LHC will provide the ideal situation for a precision equal or better than radiative corrections !
- Already the first year @ LHC will improve limits with almost an order of magnitude